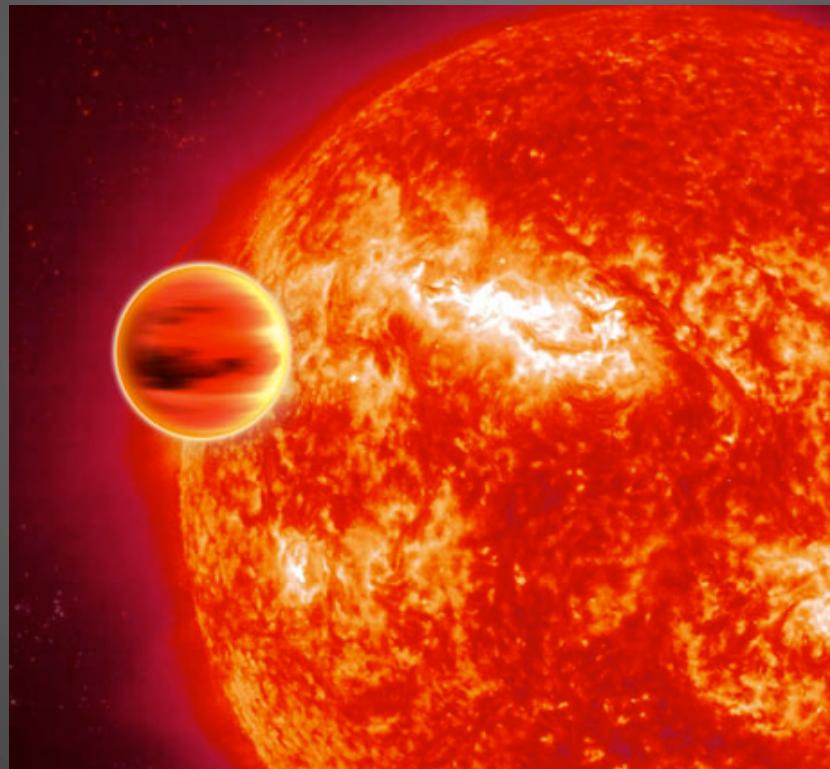


3D Models of Giant Planet Atmospheres

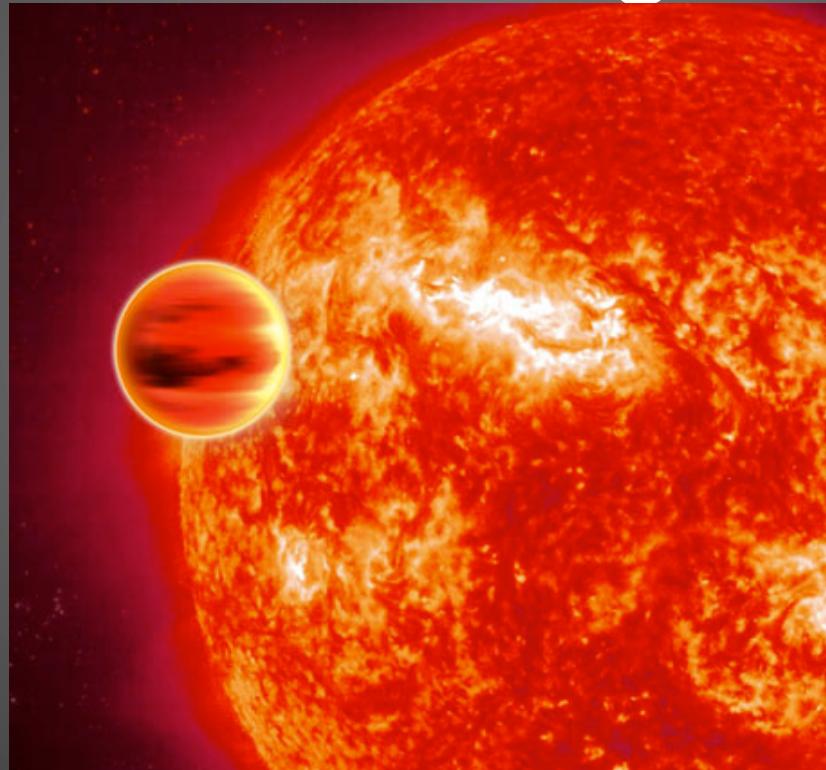


Ian Dobbs-Dixon
Sagan Postdoctoral Fellow
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Outline

- Dynamical/Radiative Modeling Methodologies
- General Properties/Thermal Inversions (x2)
- Variations of Viscosity, Changing Jet Structures
- Shocks/ Potential Vorticity Generation
- Weather
- Observable Consequences
- Eccentric Planets

Dynamical Modeling Methodologies



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Dynamical Methods

Completeness

- Equivalent Barotropic and Shallow Water (2D)
 - Cho et al (2003,2008) Langton and Laughlin (2007,2008)
Rauscher et al (2007, 2008)
- Navier-Stokes equation (2D)
 - Burkert et al. 2007
- Primitive equations (~3D)
 - Showman et al. (2002, 2005, 2006, 2008, 2009), Menou and Rauscher (2009)
- Eulers equations (3D)
 - Dobbs-Dixon and Lin (2008)
- Navier-Stokes equations (3D)
 - Dobbs-Dixon et al (2010)

Resolution

Radiation Transfer Methods

- Relaxation methods (Newtonian heating)
 - Cho et al (2003,2008) Langton and Laughlin (2007,2008) Rauscher et al (2007, 2008), Showman et al. (2002, 2005, 2006, 2008), Menou and Rauscher (2009)
- 2/3D one temperature flux-limited radiative diffusion
 - Burkert et al. (2007), Dobbs-Dixon and Lin (2008)
- 3D FLD + decoupled thermal and radiative components
 - Dobbs-Dixon et al (2009)
- 1D (radial) wavelength-dependent radiative transfer
 - Showman et al. (2009)

East-West

$$\rho \frac{\partial u}{\partial t} + u \frac{\partial \rho}{\partial t} = - \frac{1}{r \cos \phi} \frac{\partial p}{\partial \lambda} + \rho \left(2\Omega + \frac{u}{r \cos \phi} \right) (\sin \phi) v - \rho \left(2\Omega \cos \phi + \frac{u}{r} \right) w \\ - \left[\frac{1}{r \cos \phi} \frac{\partial}{\partial \lambda} (\rho u^2) + \frac{1}{r \cos \phi} \frac{\partial}{\partial \phi} (\rho uv \cos \phi) + \frac{1}{r^2} \frac{\partial}{\partial r} (\rho uw r^2) \right] + \rho F_\lambda,$$

North-South

$$\rho \frac{\partial v}{\partial t} + v \frac{\partial \rho}{\partial t} = - \frac{1}{r} \frac{\partial p}{\partial \phi} - \rho \left(2\Omega + \frac{u}{r \cos \phi} \right) (\sin \phi) u - \rho \frac{vw}{r} \\ - \left[\frac{1}{r \cos \phi} \frac{\partial}{\partial \lambda} (\rho uv) + \frac{1}{r \cos \phi} \frac{\partial}{\partial \phi} (\rho v^2 \cos \phi) + \frac{1}{r^2} \frac{\partial}{\partial r} (\rho r^2 vw) \right] + \rho F_\phi,$$

Radial

$$\rho \frac{\partial w}{\partial t} + w \frac{\partial \rho}{\partial t} = - \frac{\partial p}{\partial r} - \rho g + \rho \left(2\Omega \cos \phi + \frac{u}{r} \right) u + \rho \frac{v^2}{r} \\ - \left[\frac{1}{r \cos \phi} \frac{\partial}{\partial \lambda} (\rho uw) + \frac{1}{r \cos \phi} \frac{\partial}{\partial \phi} (\rho vw \cos \phi) + \frac{1}{r^2} \frac{\partial}{\partial r} (\rho r^2 w^2) \right] + \rho F_r.$$

Gilman and Glatzmaier 1981

3D Navier-Stokes, flux limited diffusion and decoupled thermal and radiative components

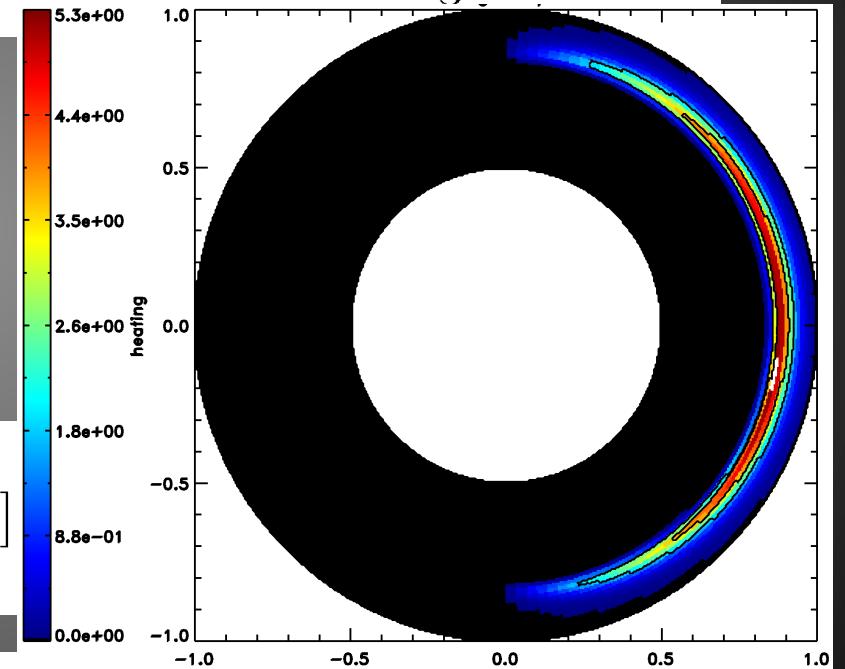
$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{\nabla P}{\rho} + \mathbf{g} - 2\boldsymbol{\Omega} \times \mathbf{u} - \boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{r}) + \nu \nabla^2 \mathbf{u} + \frac{\nu}{3} \nabla (\nabla \cdot \mathbf{u})$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

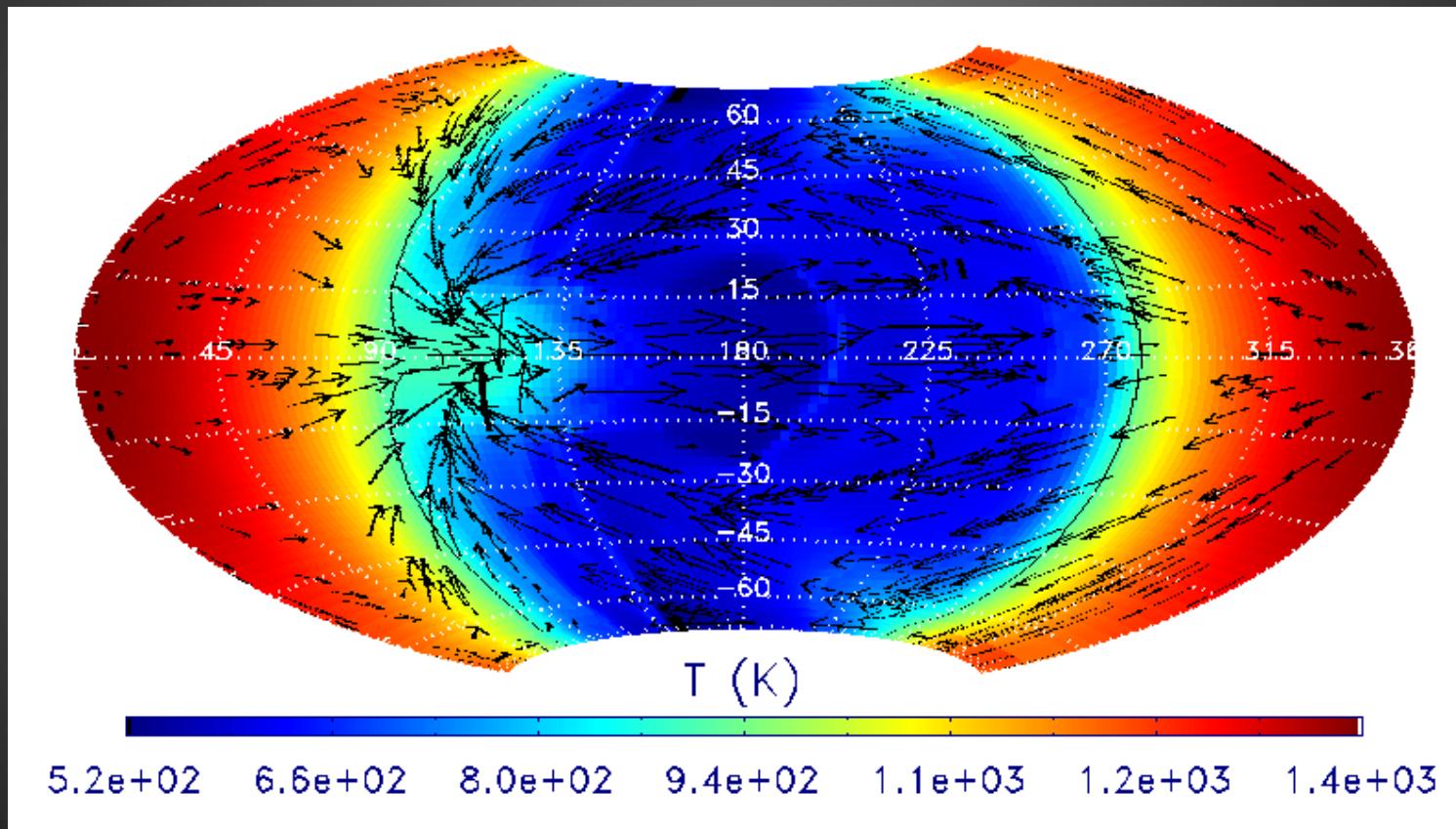
$$\mathbf{F} = -\lambda \frac{c}{\rho \kappa_R(T, P)} \nabla E_R$$

$$\frac{\partial E_R}{\partial t} + \nabla \cdot \mathbf{F} = \rho \kappa_P(T, P) [B(T) - cE_R]$$

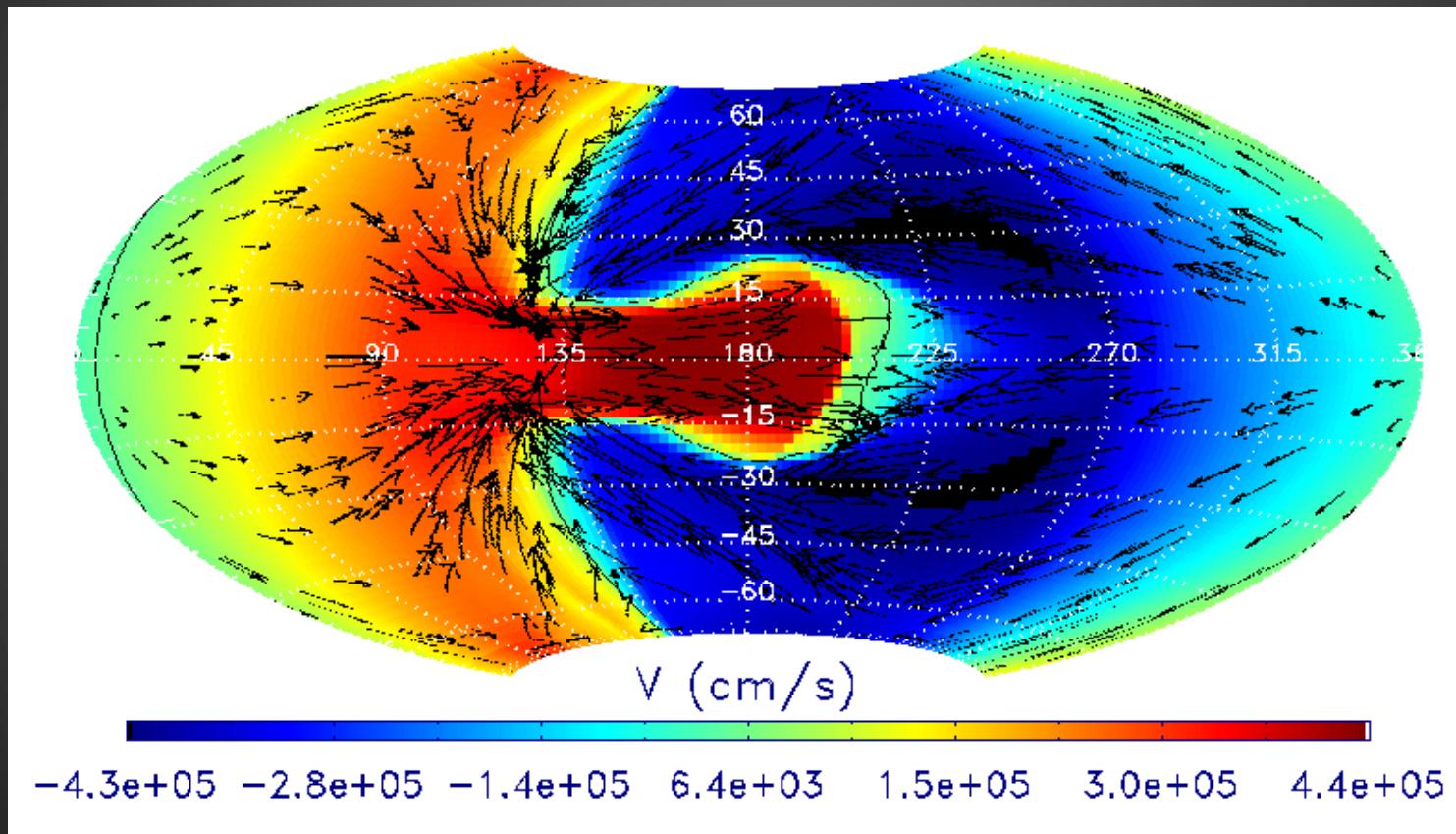
$$\left[\frac{\partial \epsilon}{\partial t} + (\mathbf{u} \cdot \nabla) \epsilon \right] = -P \nabla \cdot \mathbf{u} - \mu \kappa_P(T, P) [B(T) - cE_R] + \rho \kappa_{\star}(T, P) F_{\star} e^{-\tau_{\star}} + \Phi_{\nu}.$$



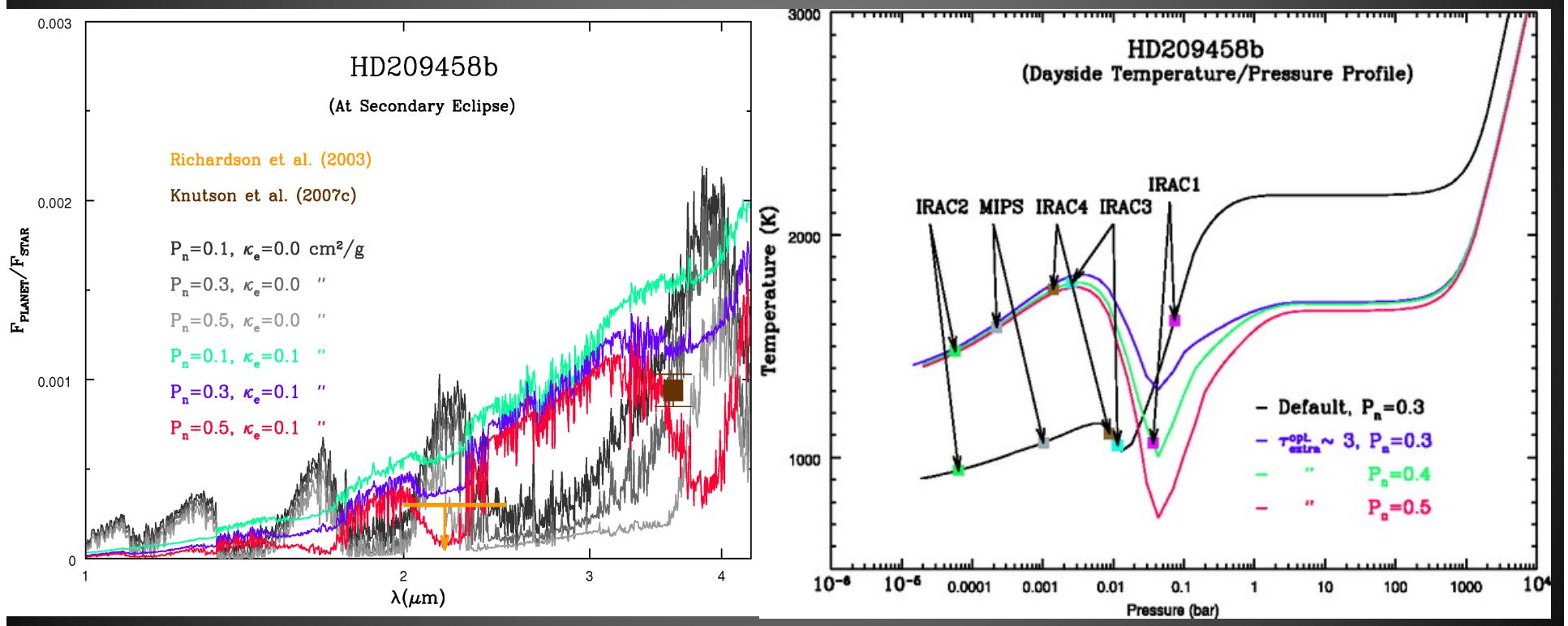
$P_{\text{rot}} = P_{\text{orb}} = 3.52\text{d}$, $T_{\star} = 6117\text{K}$
 $M_p = 0.69M_J$, $R_p = 1.3R_J$



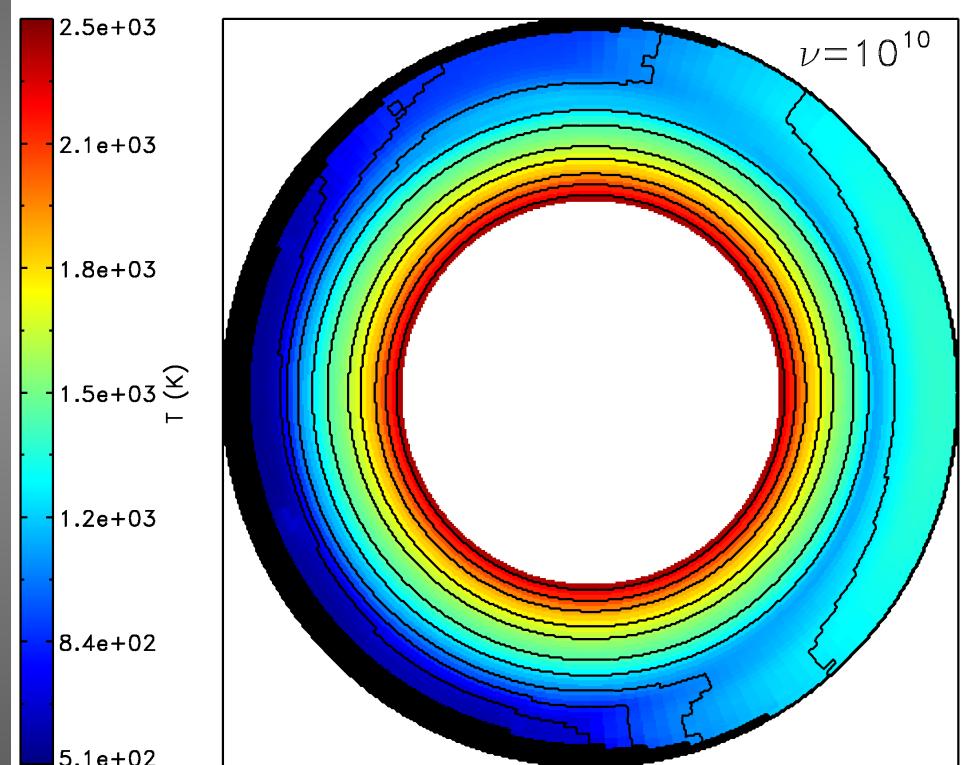
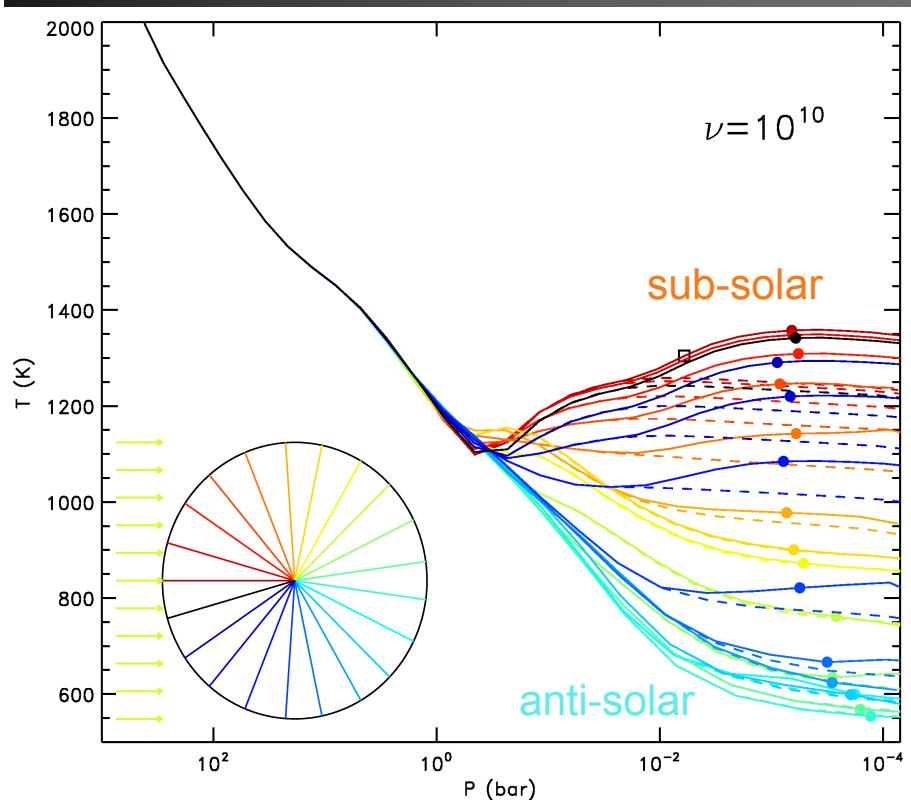
Photospheric Velocities



Observed Inversion (HD 209458b)

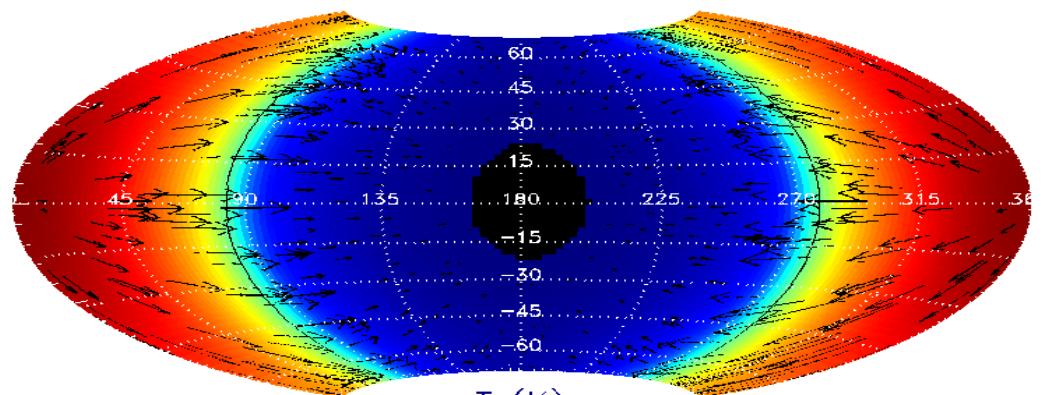


HD209458b

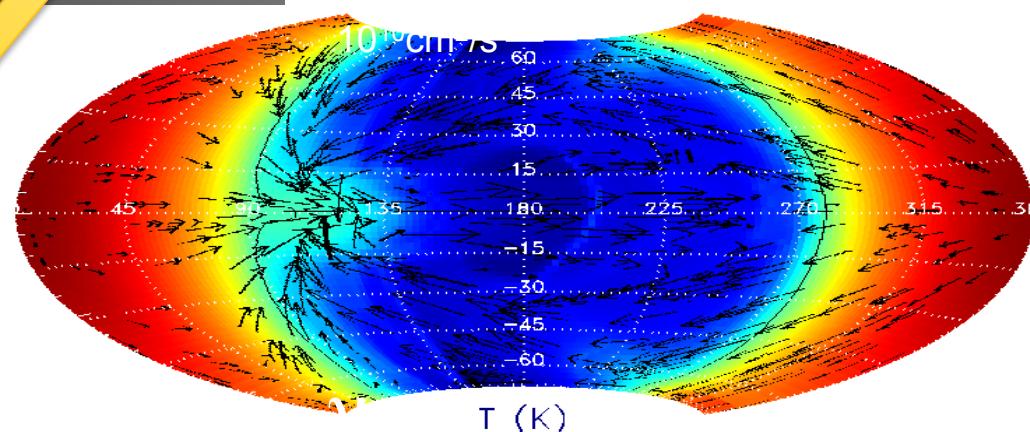


Varying Viscosity

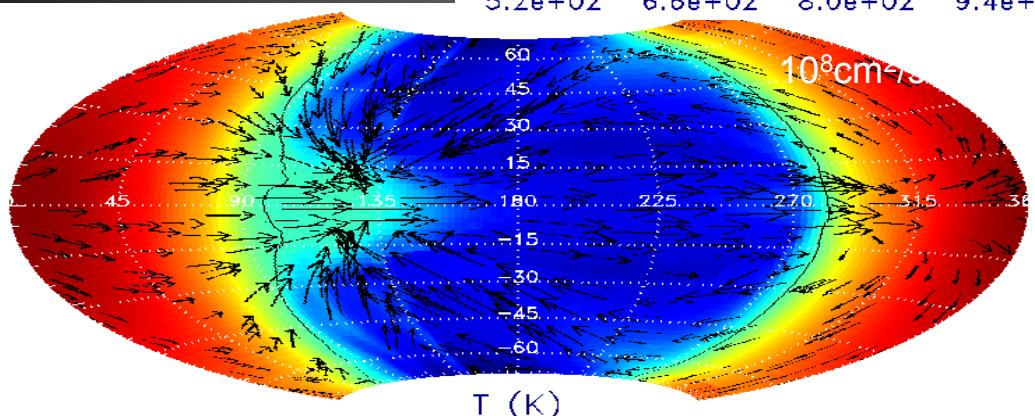
Viscosity ↑



$10^6 \text{ cm}^2/\text{s}$



$10^7 \text{ cm}^2/\text{s}$

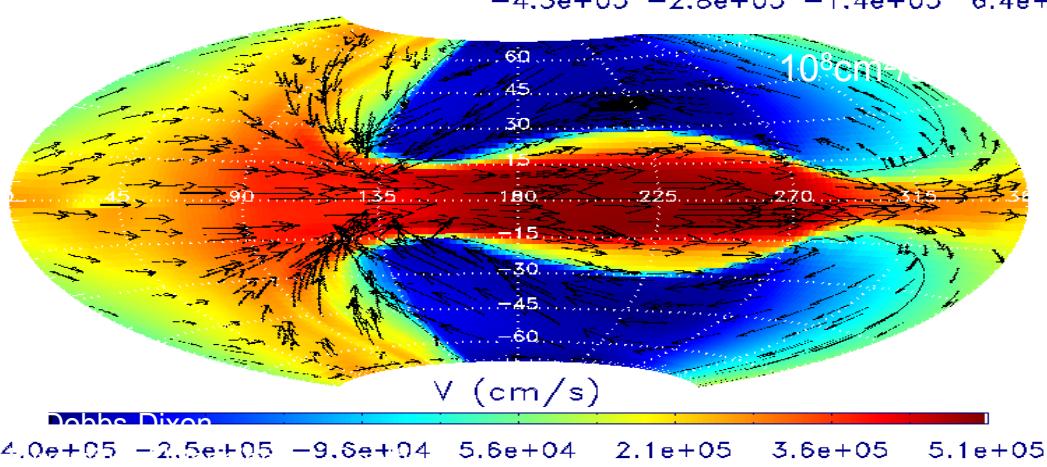
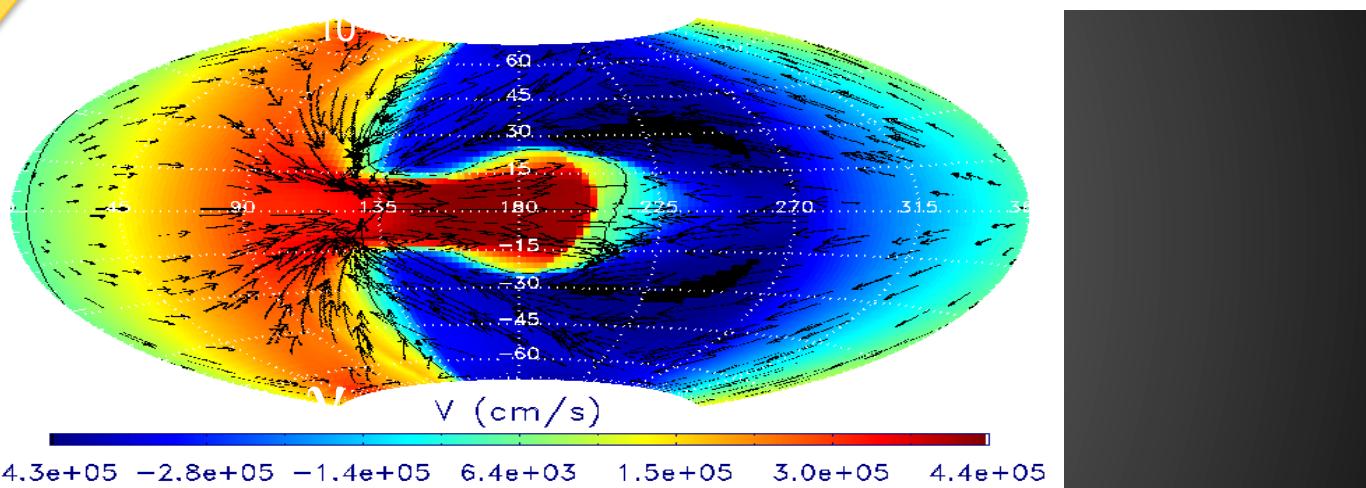
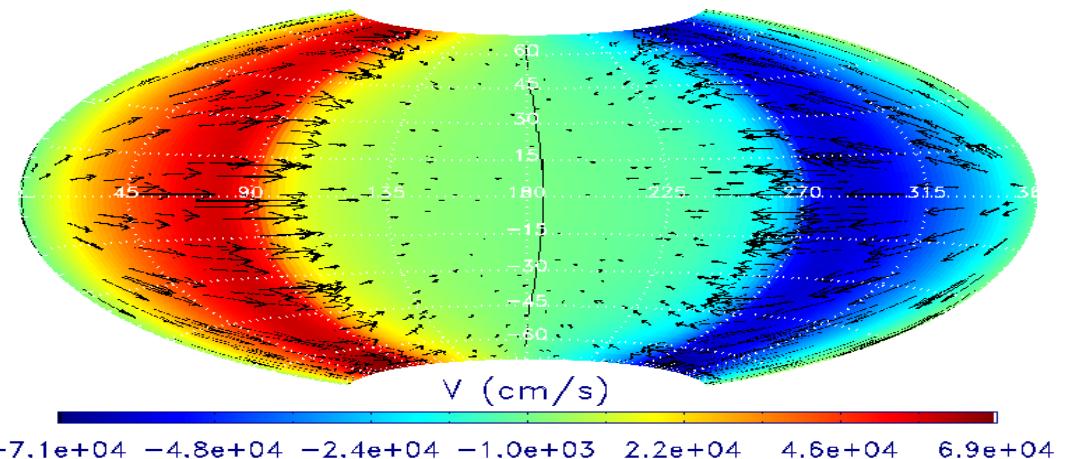


$10^8 \text{ cm}^2/\text{s}$

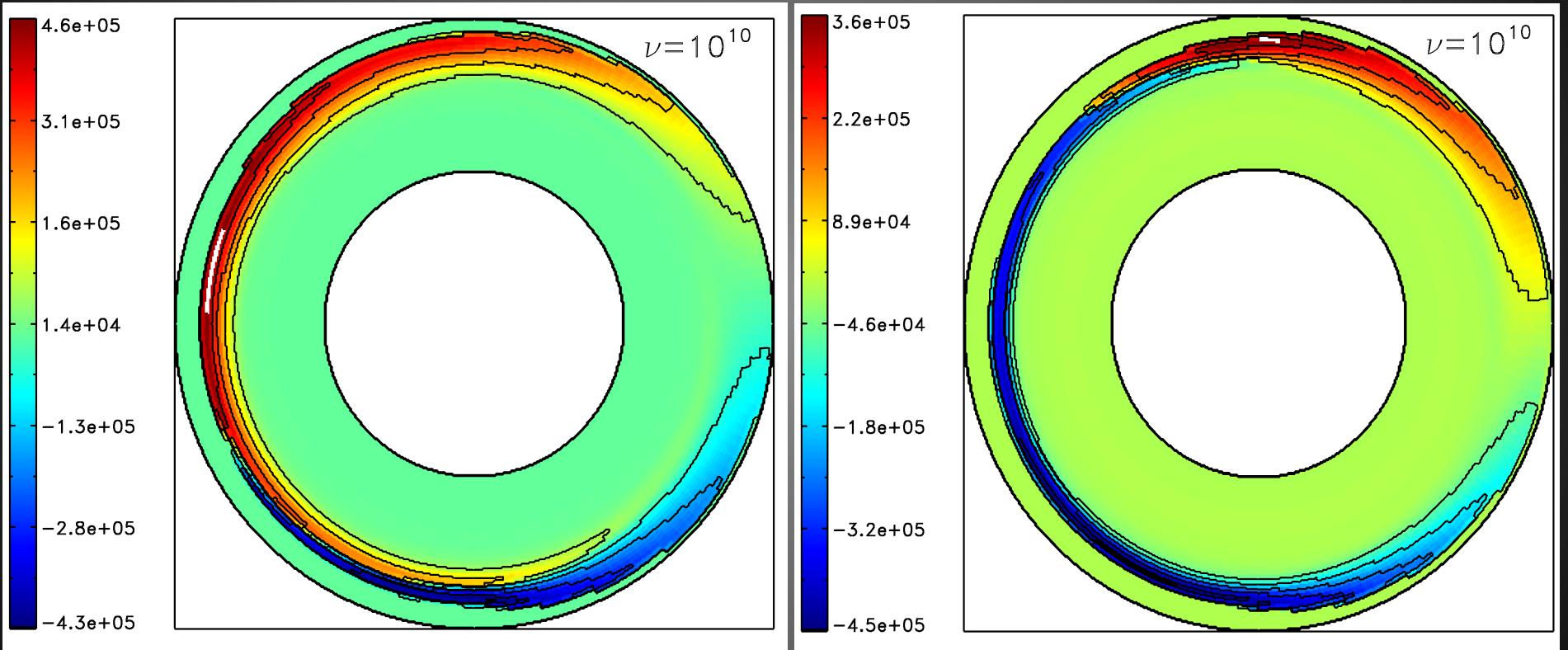


Varying Viscosity

Viscosity ↑



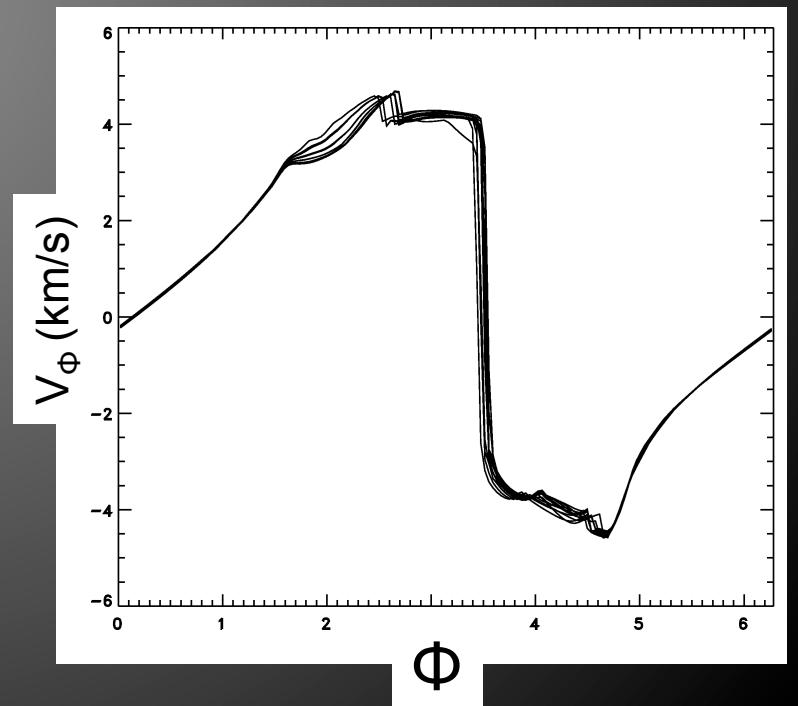
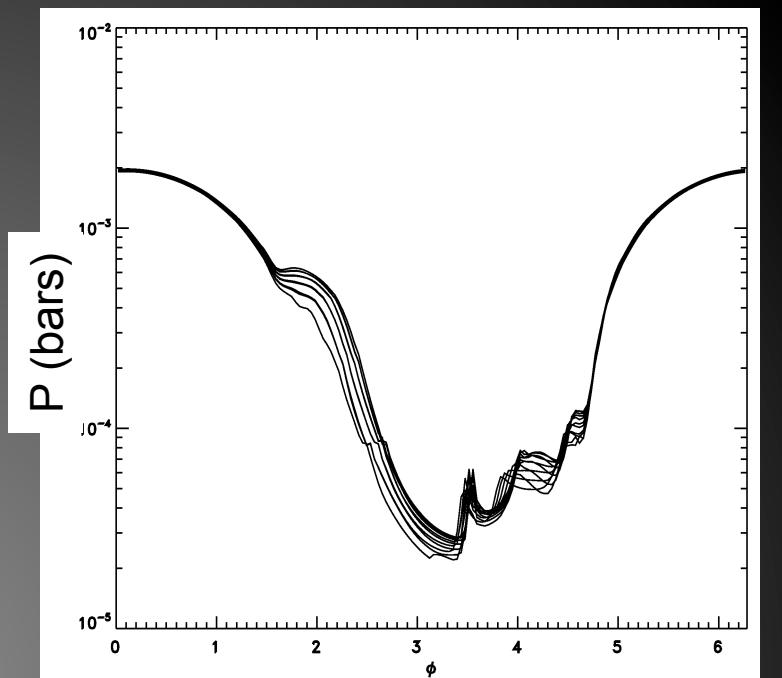
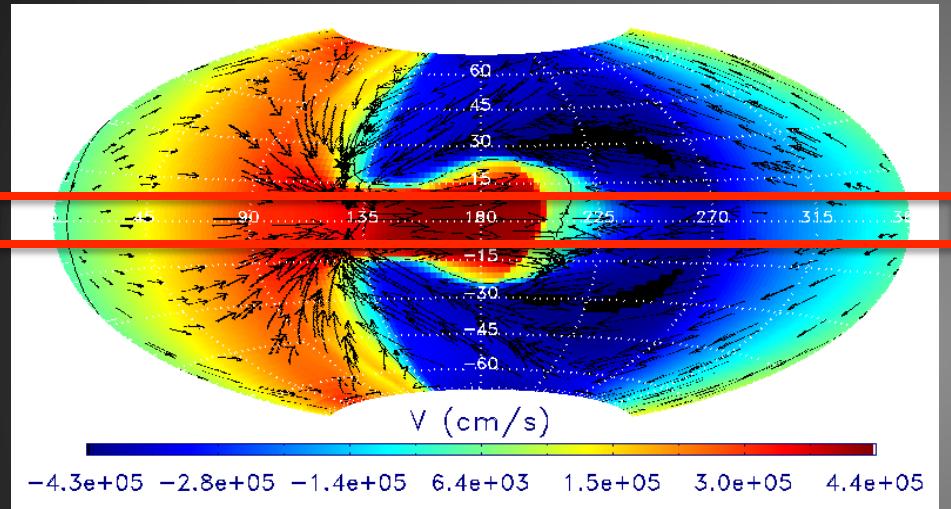
Velocity Structure with Radius



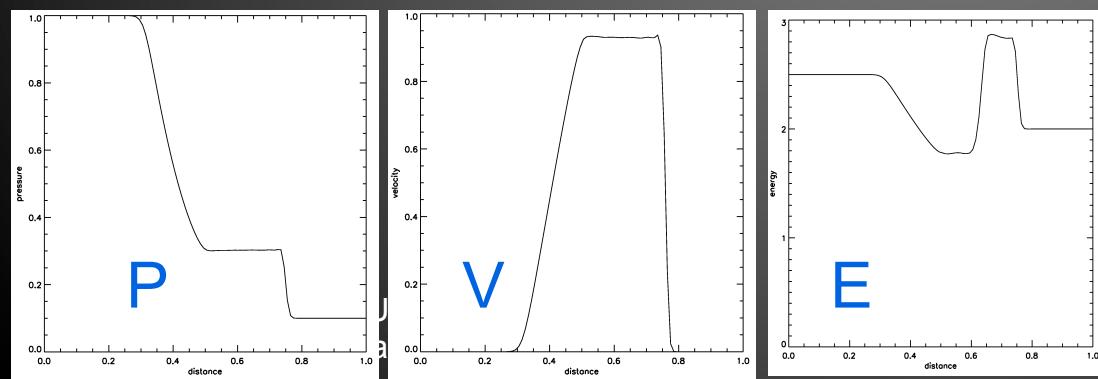
Equator

Mid-Latitude

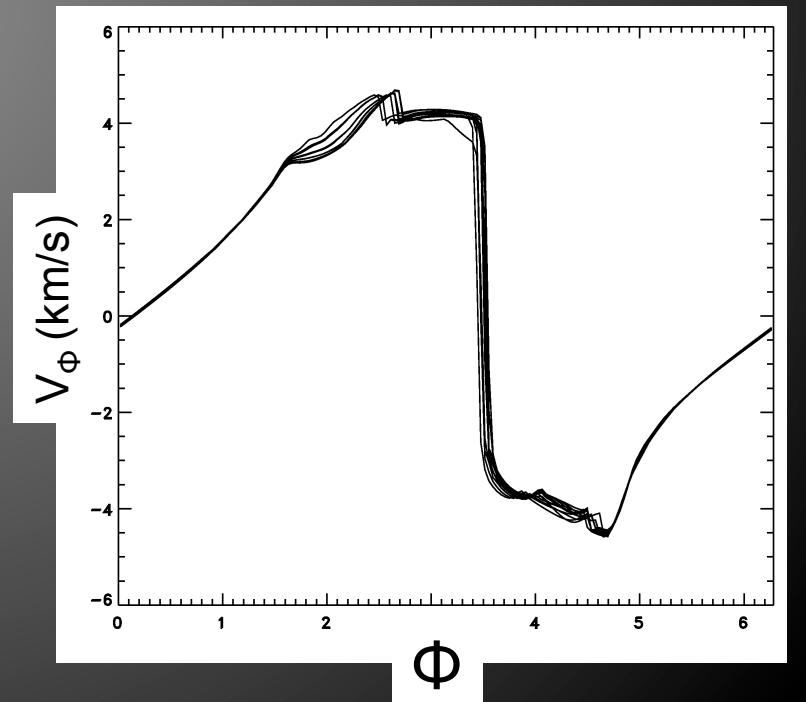
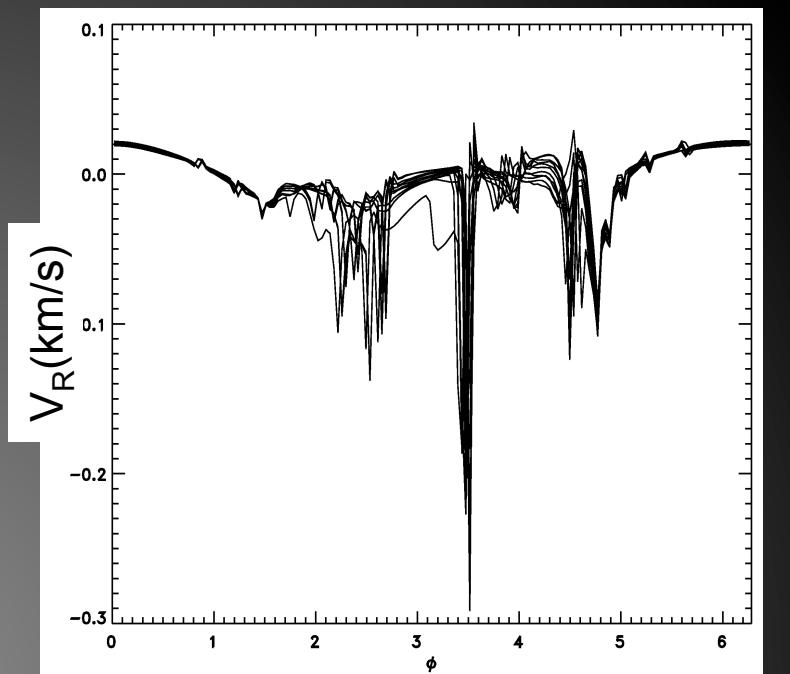
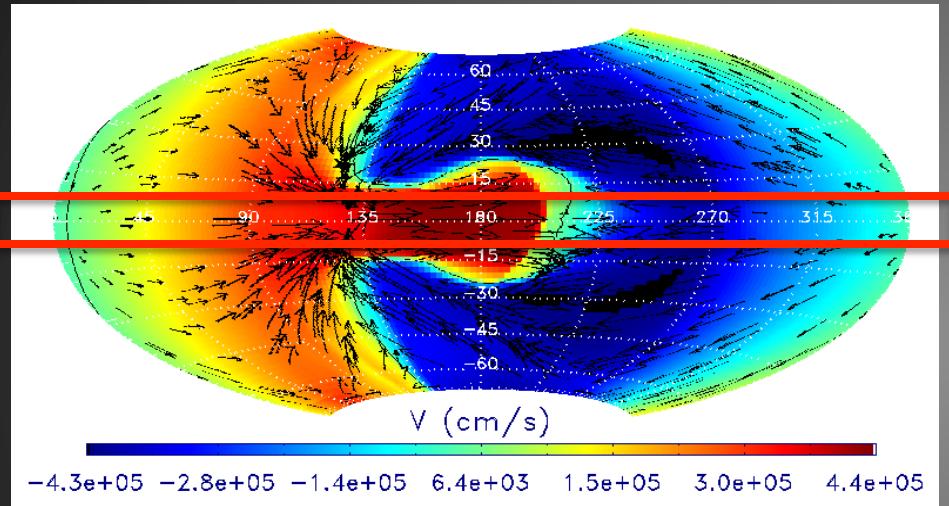
Horizontal Equatorial Shock Structure



1D Riemann shock tube

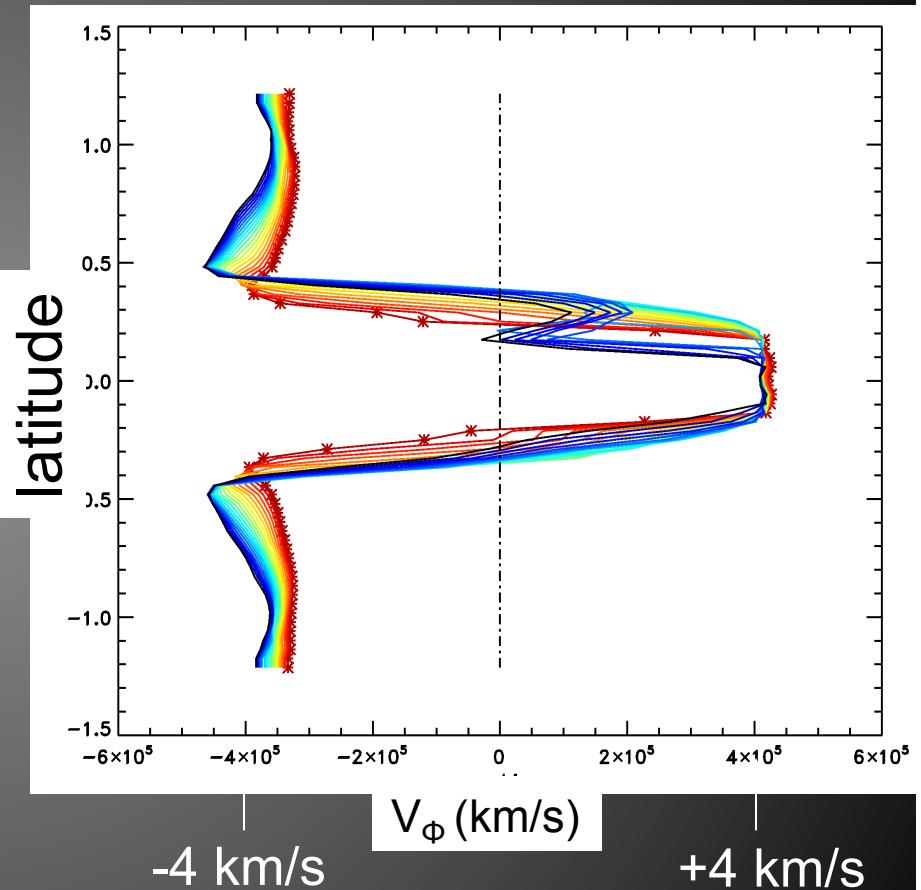
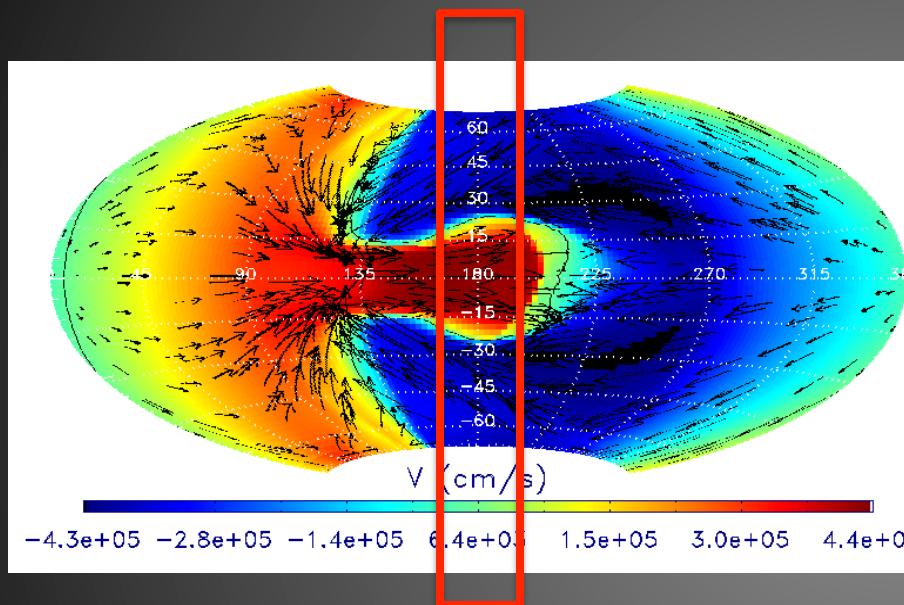


Horizontal Equatorial Shock Structure

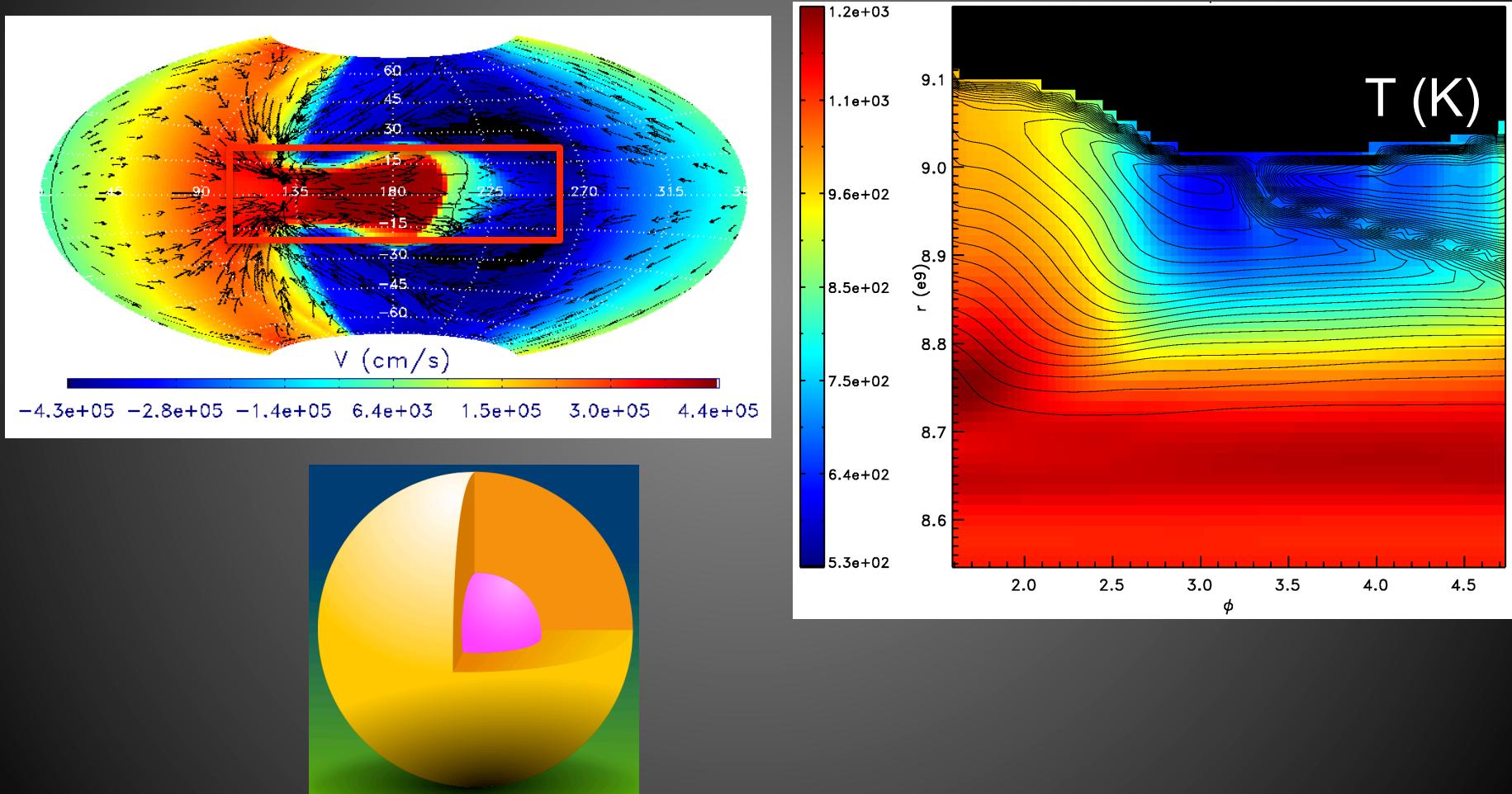


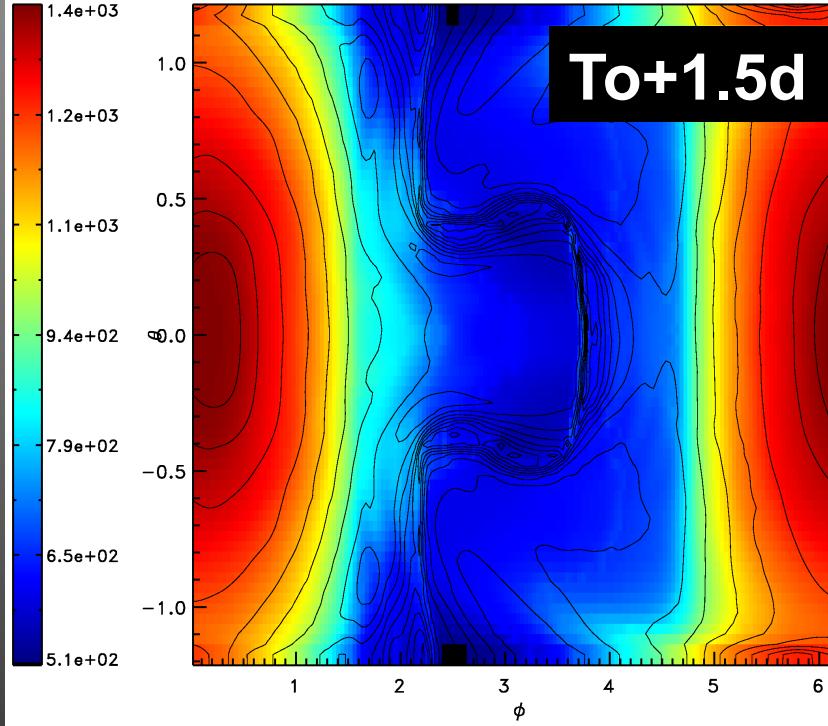
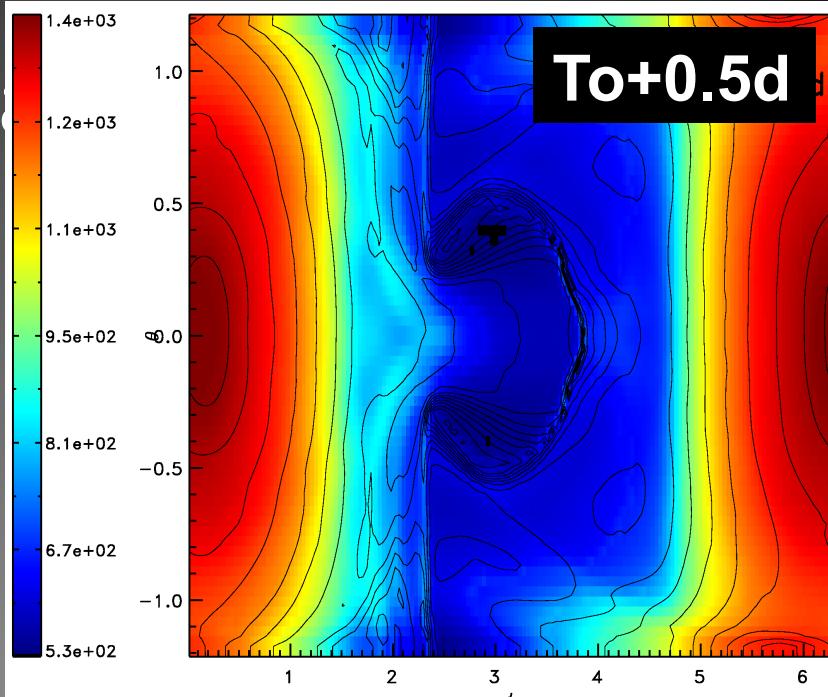
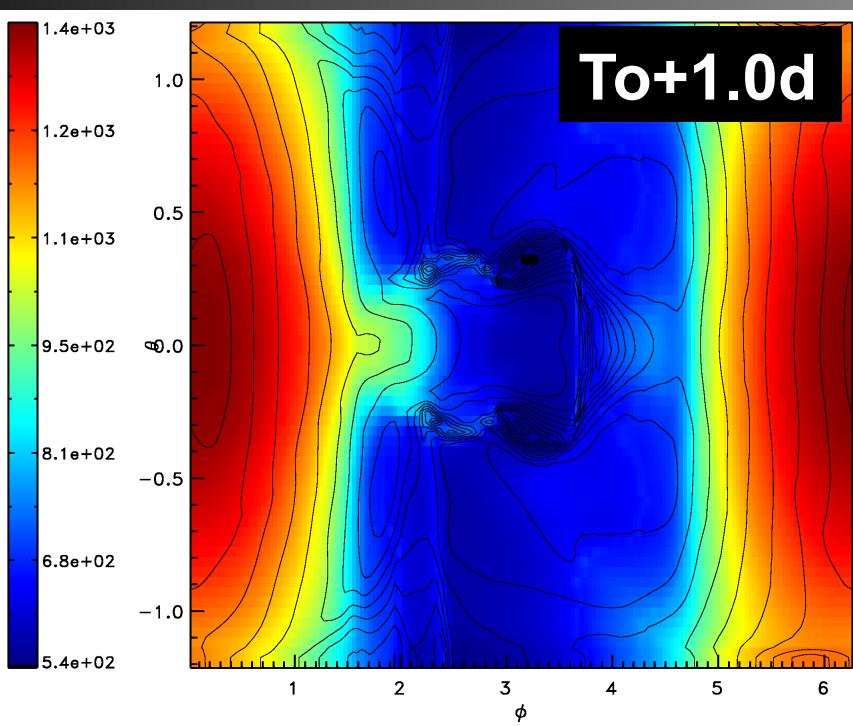
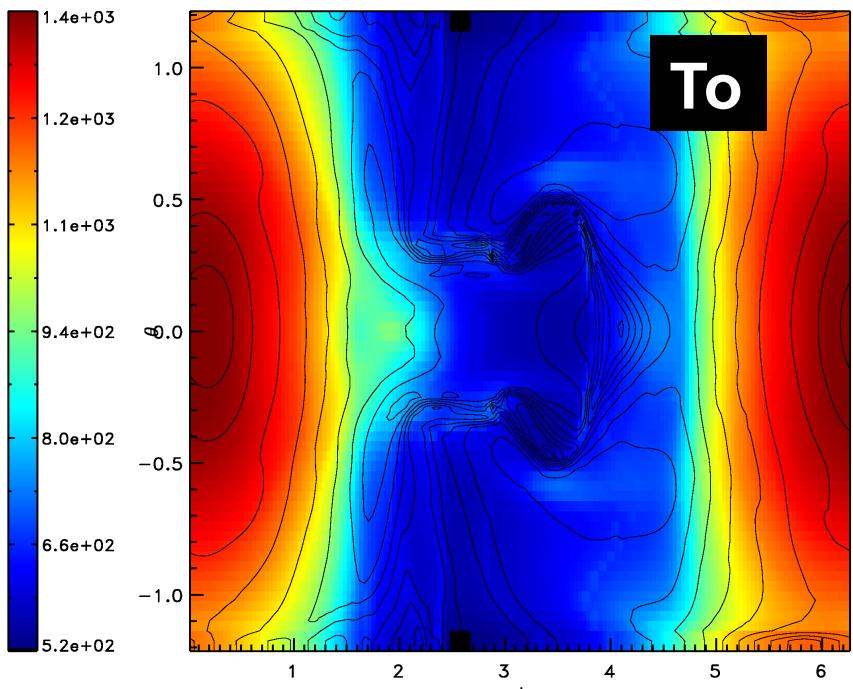
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Horizontal Latitudinal Shearing

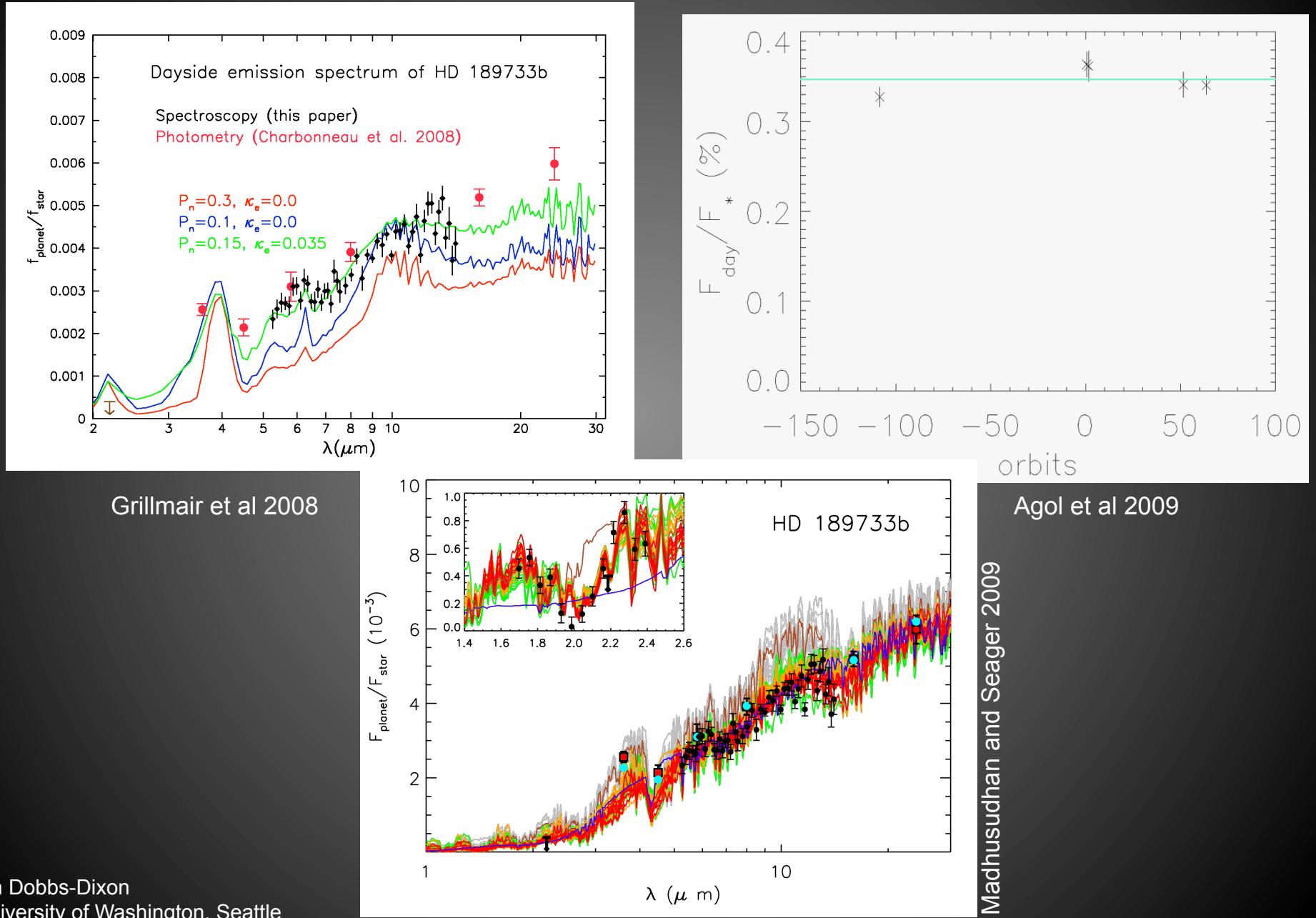


Radial Shearing

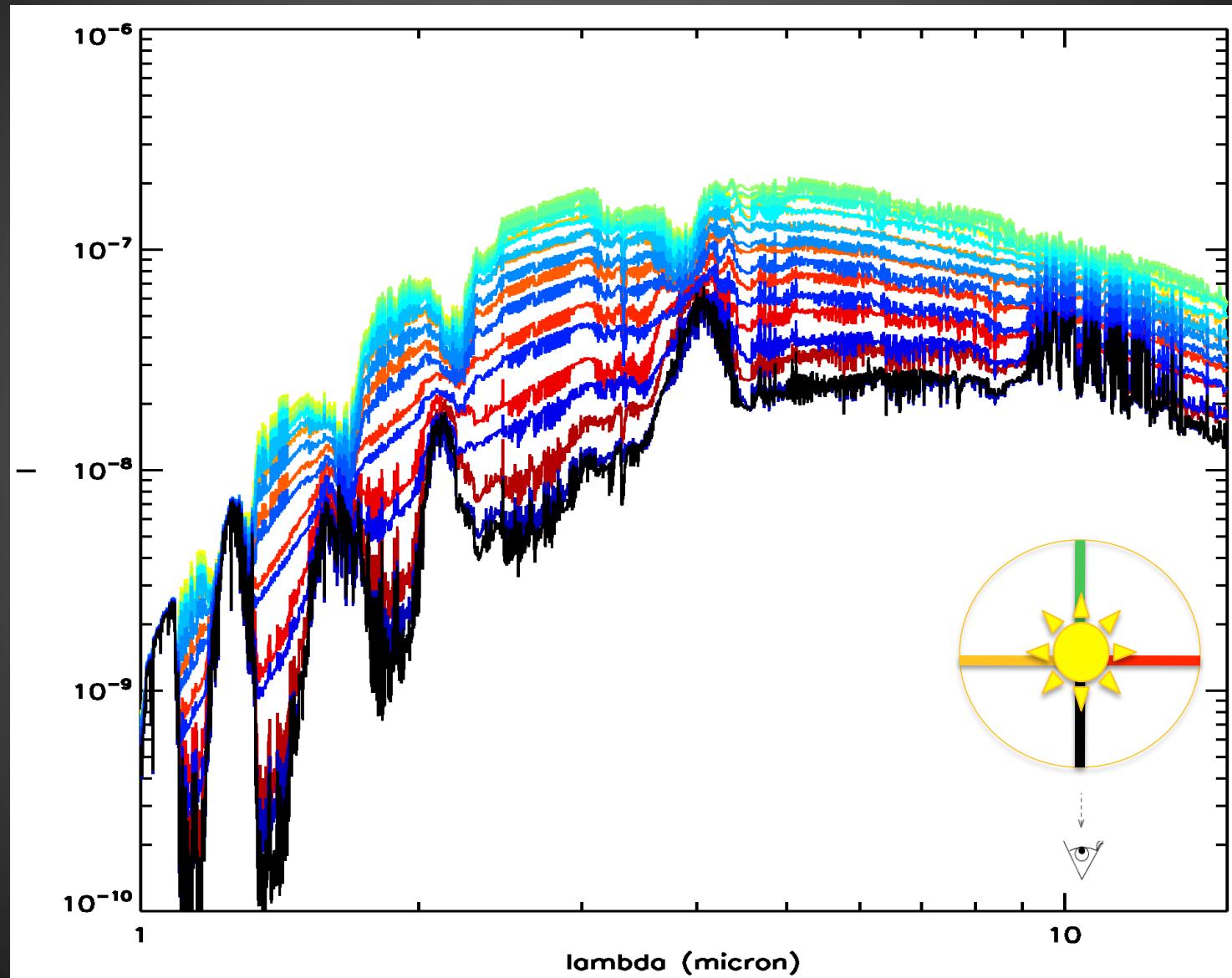




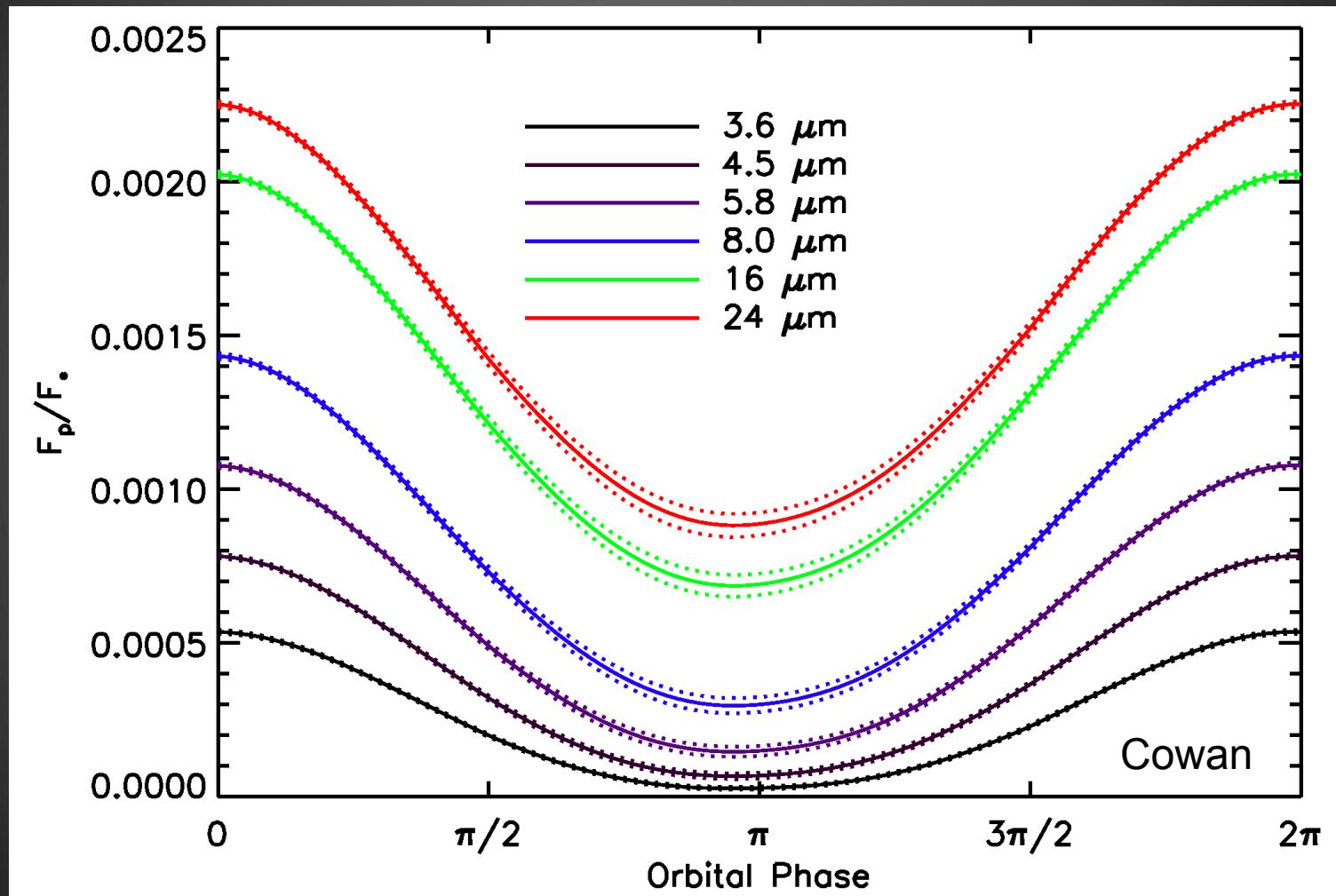
Observed Variability?



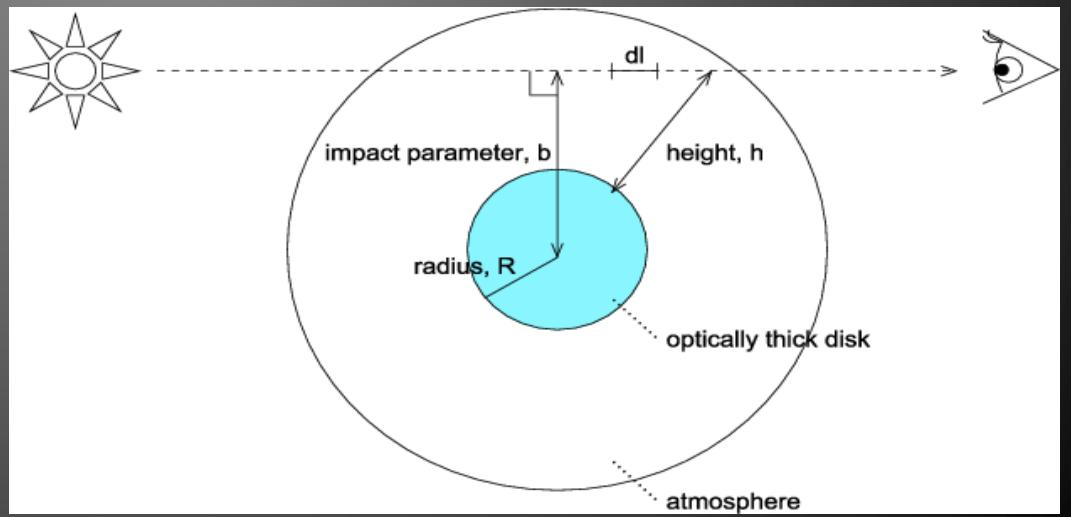
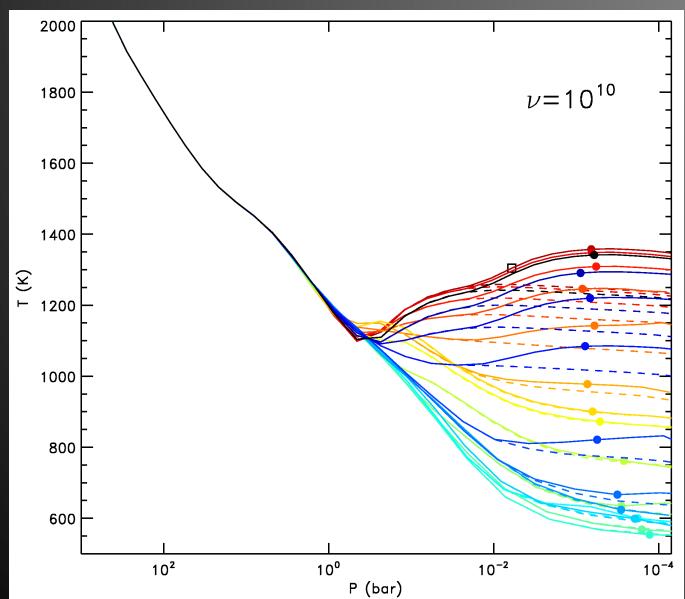
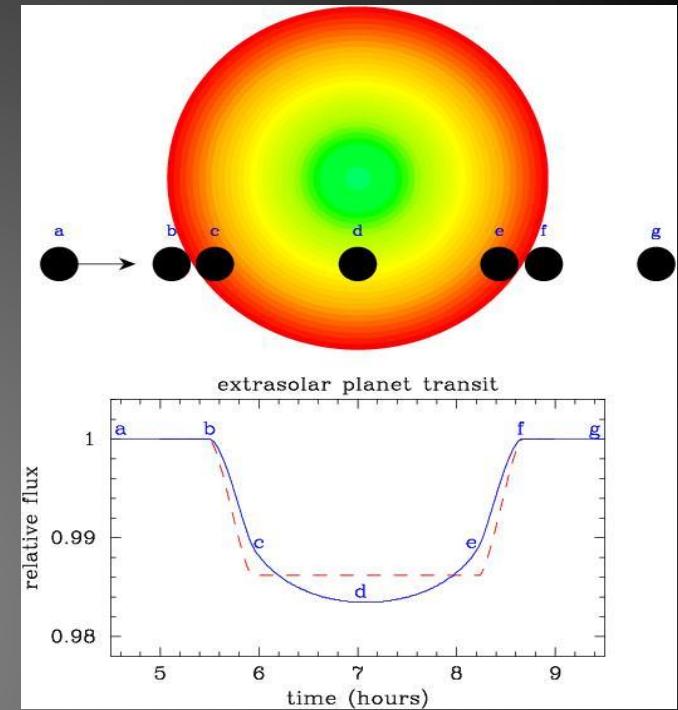
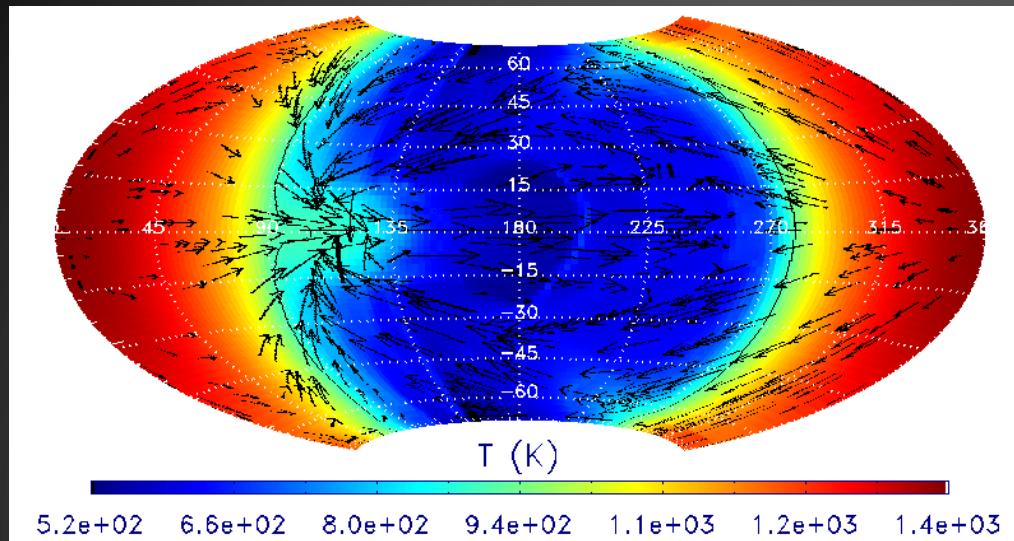
Emission Spectra



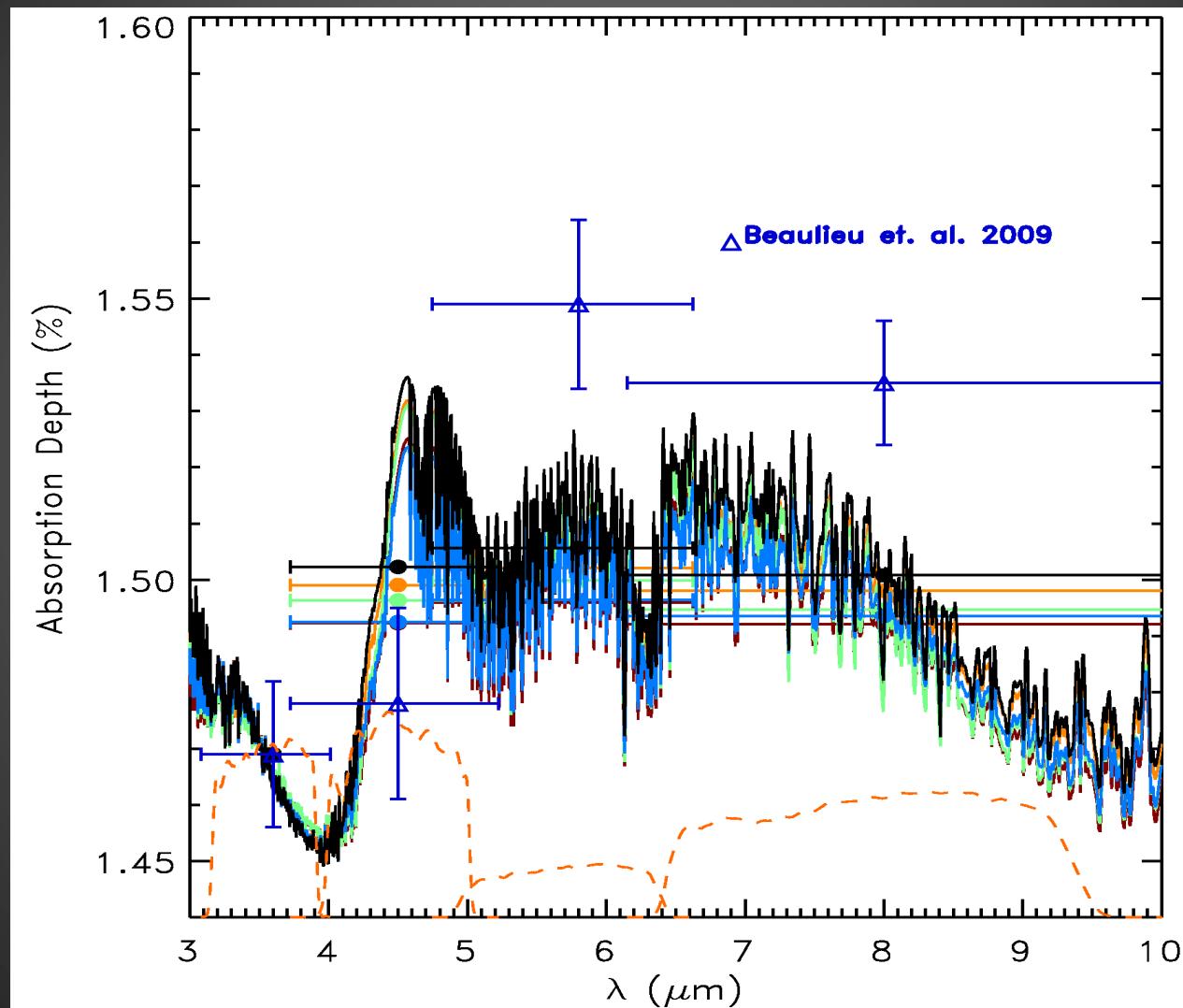
Emission Spectra



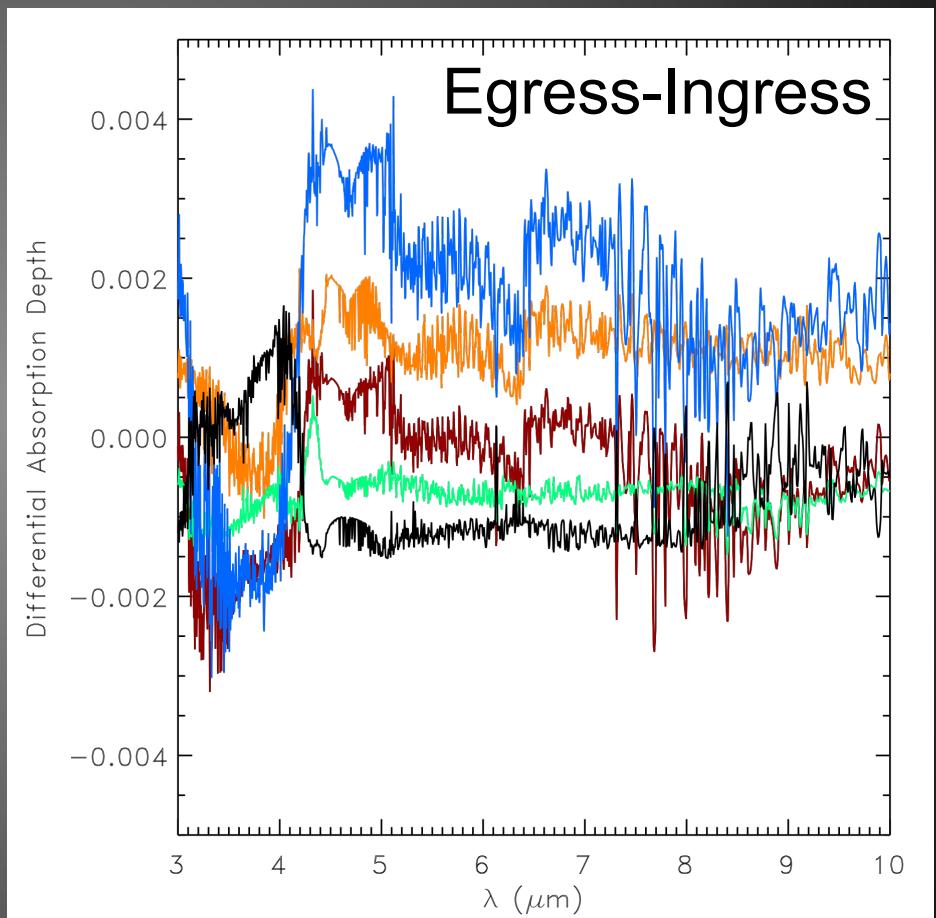
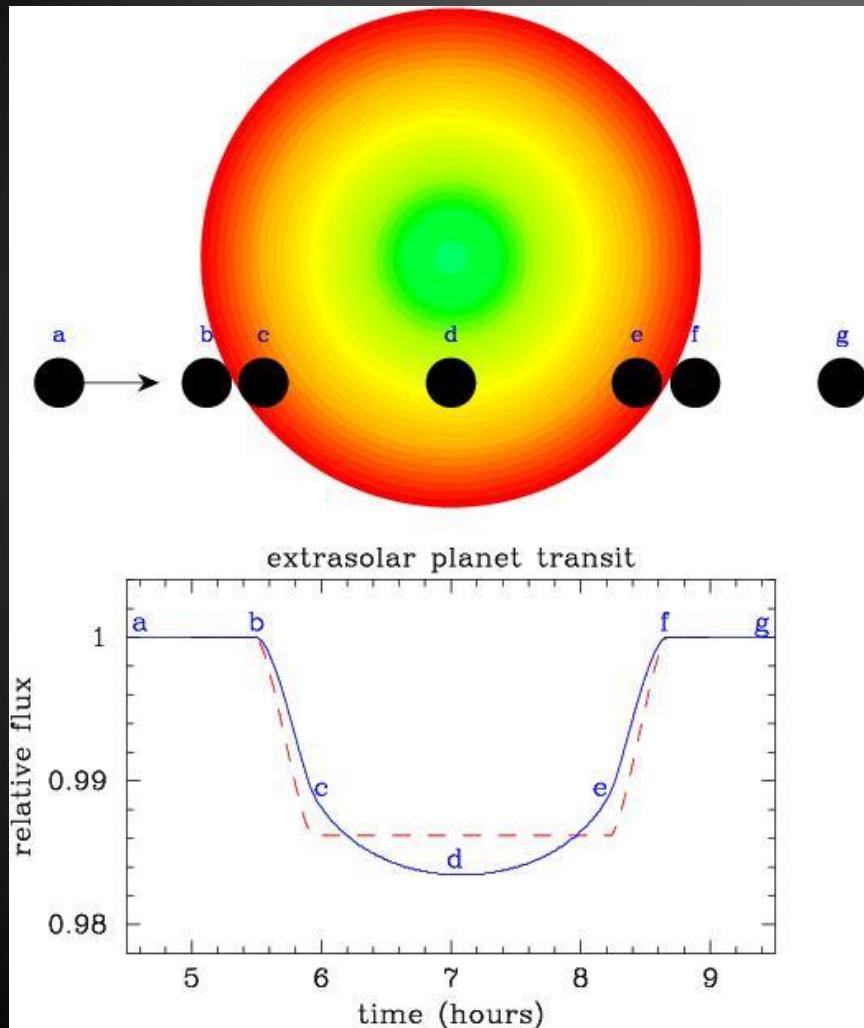
Transmission Spectra



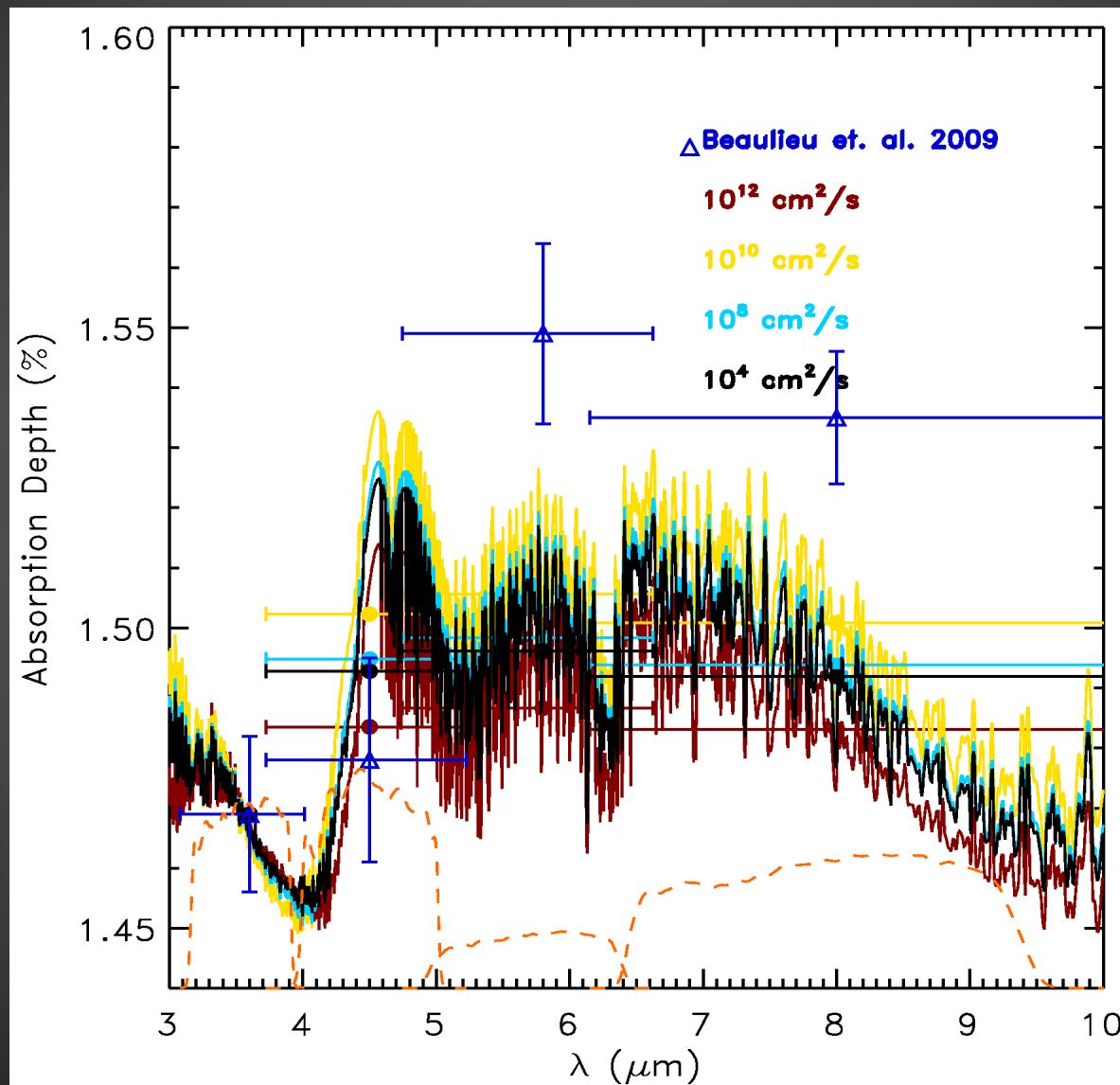
Transmission Spectra: temporal variations



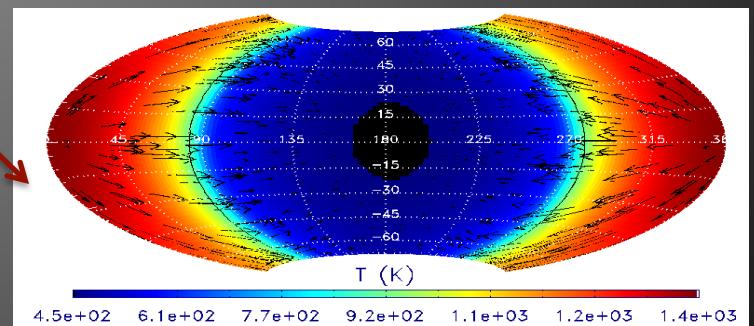
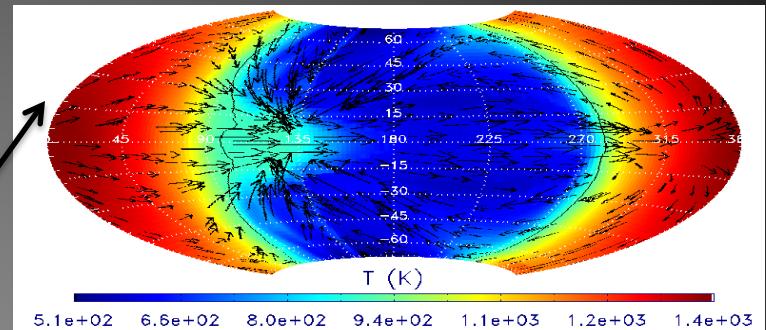
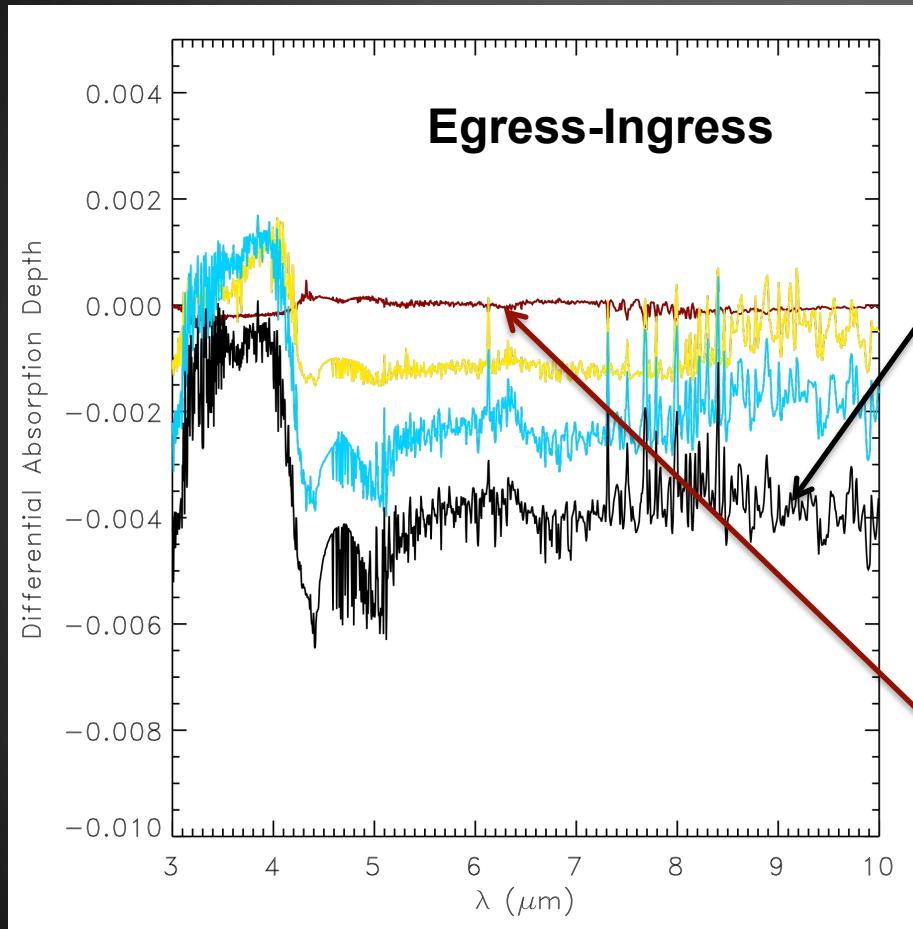
Transmission Spectra: temporal variations



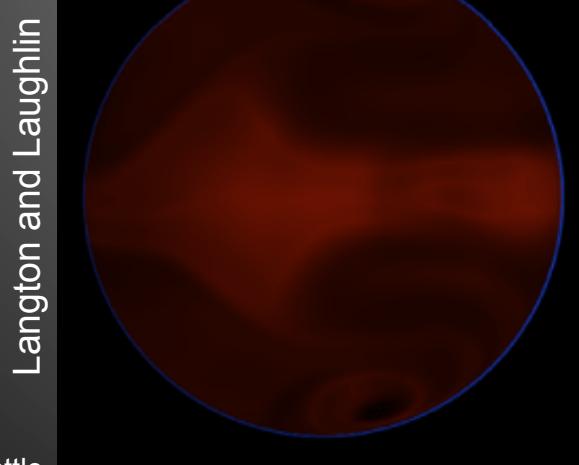
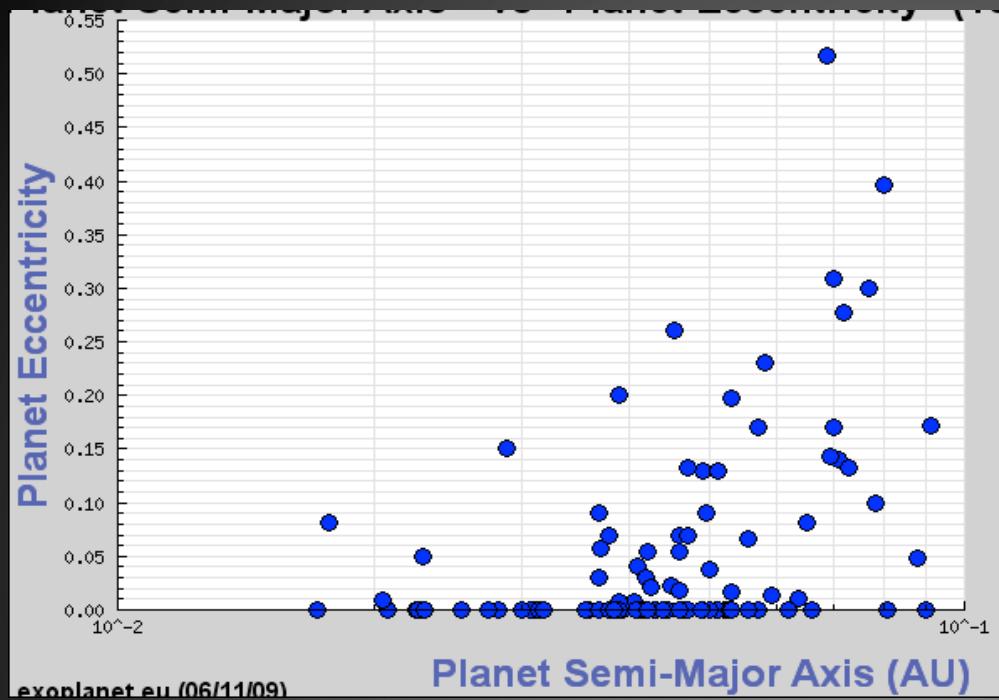
Transmission Spectra: viscous effects



Transmission Spectra: Viscous Variations



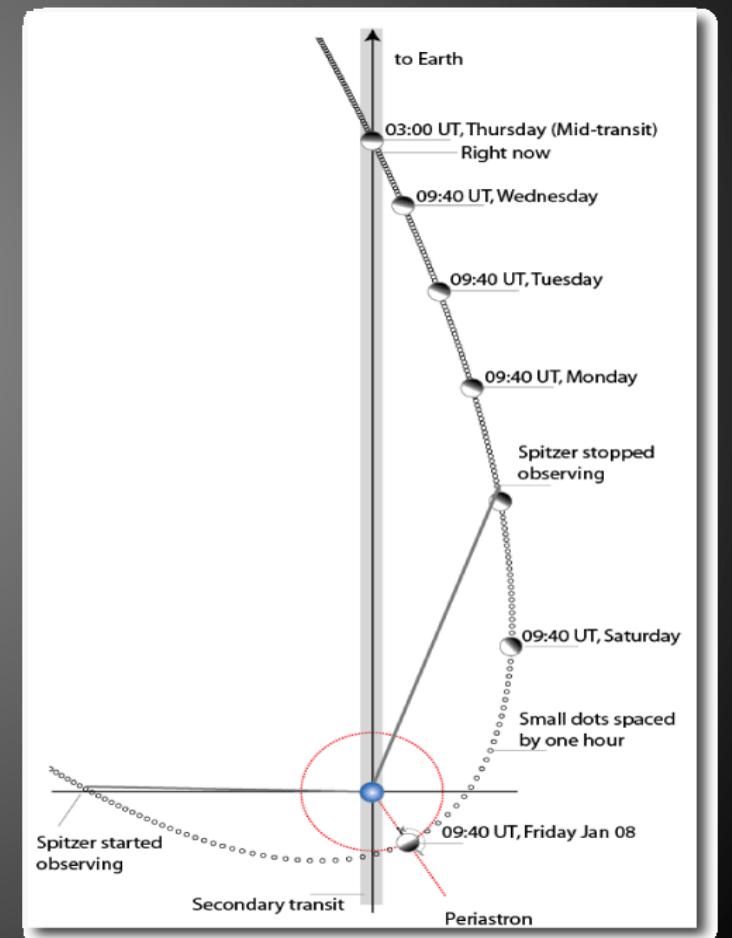
Eccentric Planets



Langton and Laughlin

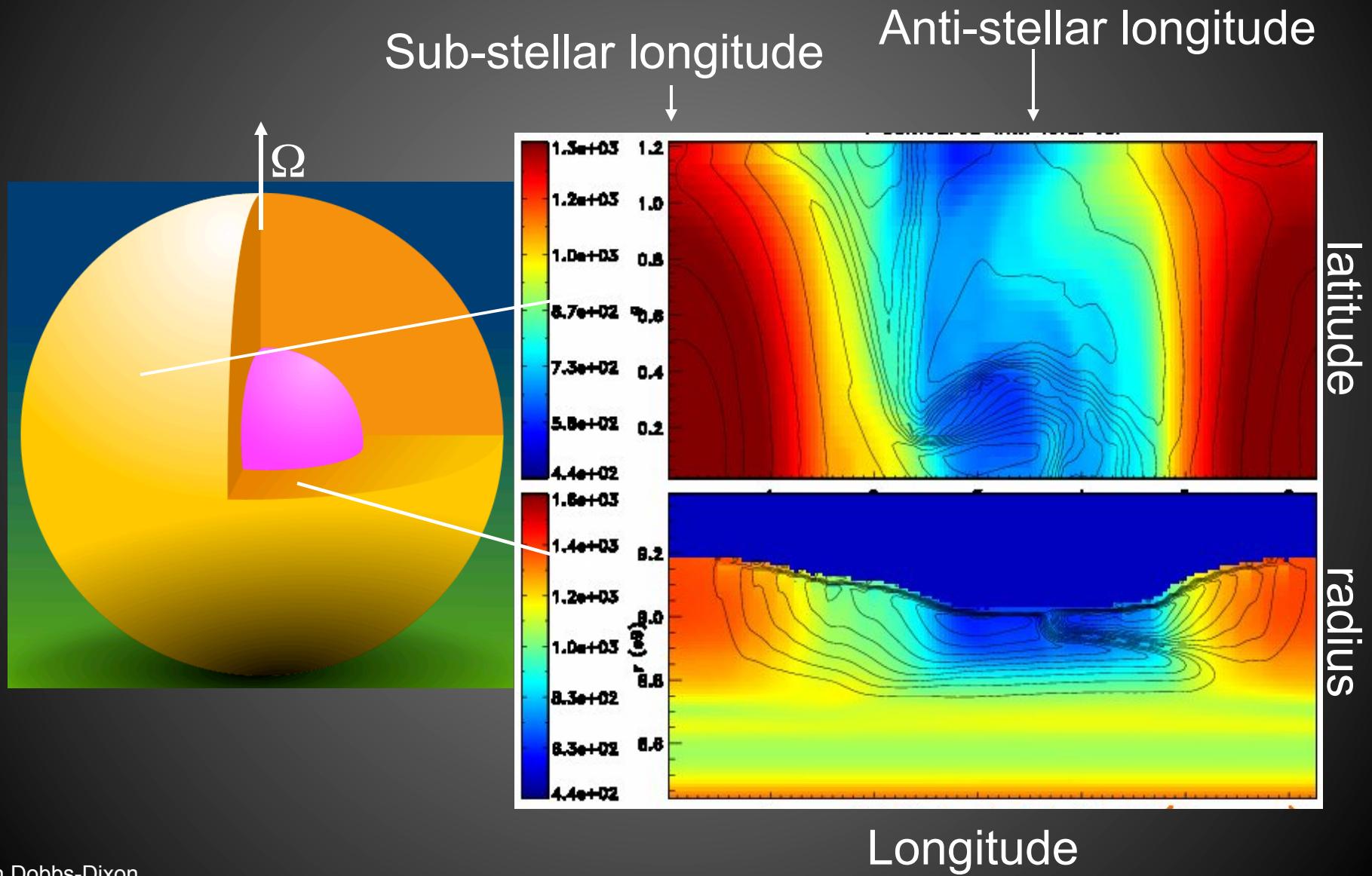
Ian Dobbs-Dixon
University of Washington, Seattle

HD80606b
 $P=111\text{days}$, $e=0.9$

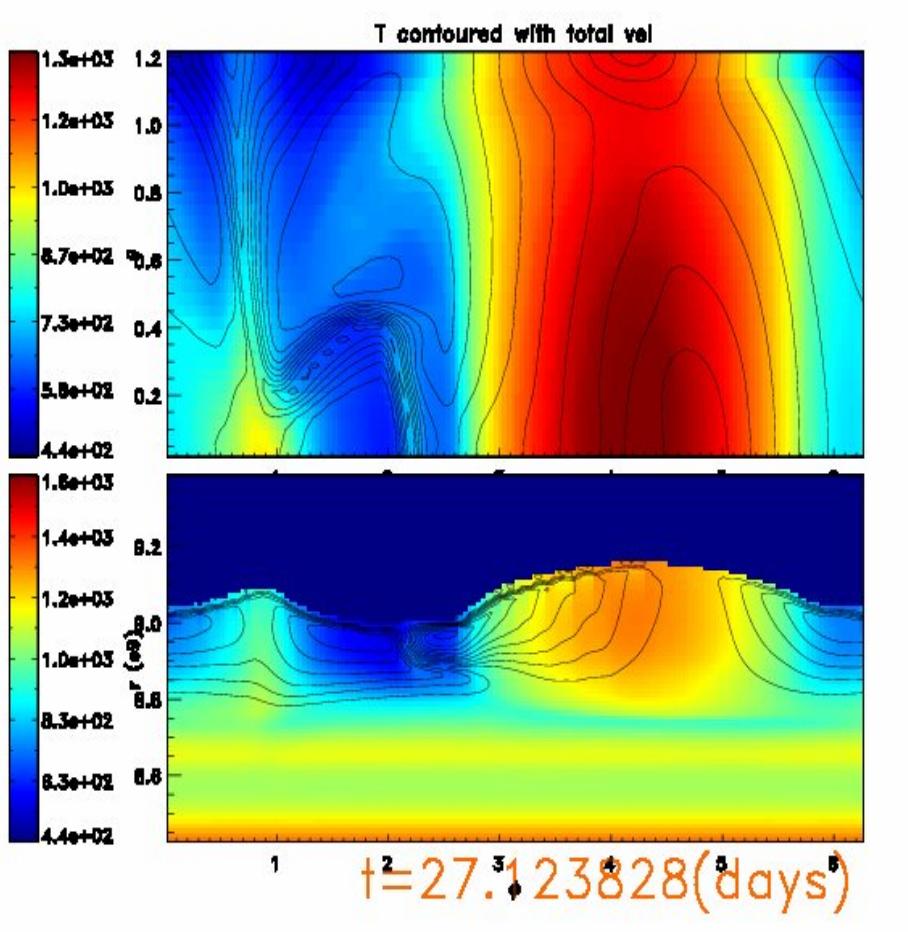


G. Laughlin

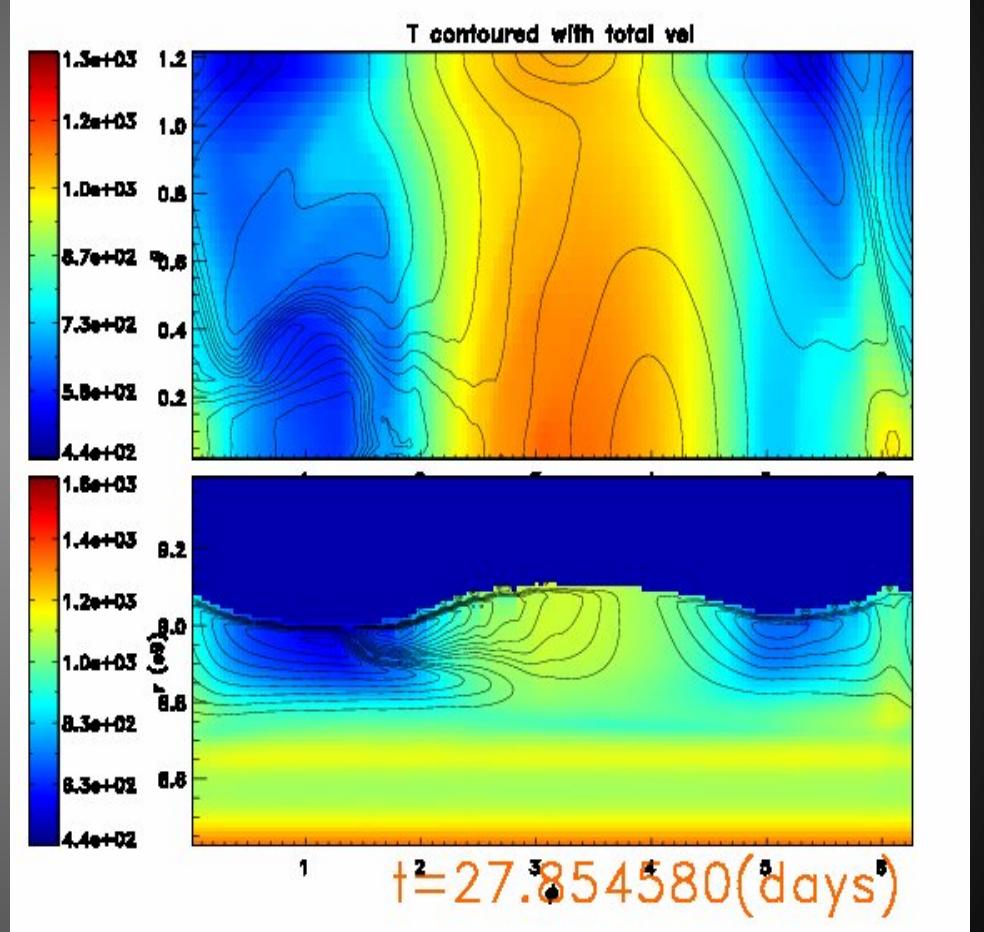
Eccentric Planets



Eccentric Planets ($e=0.5$)

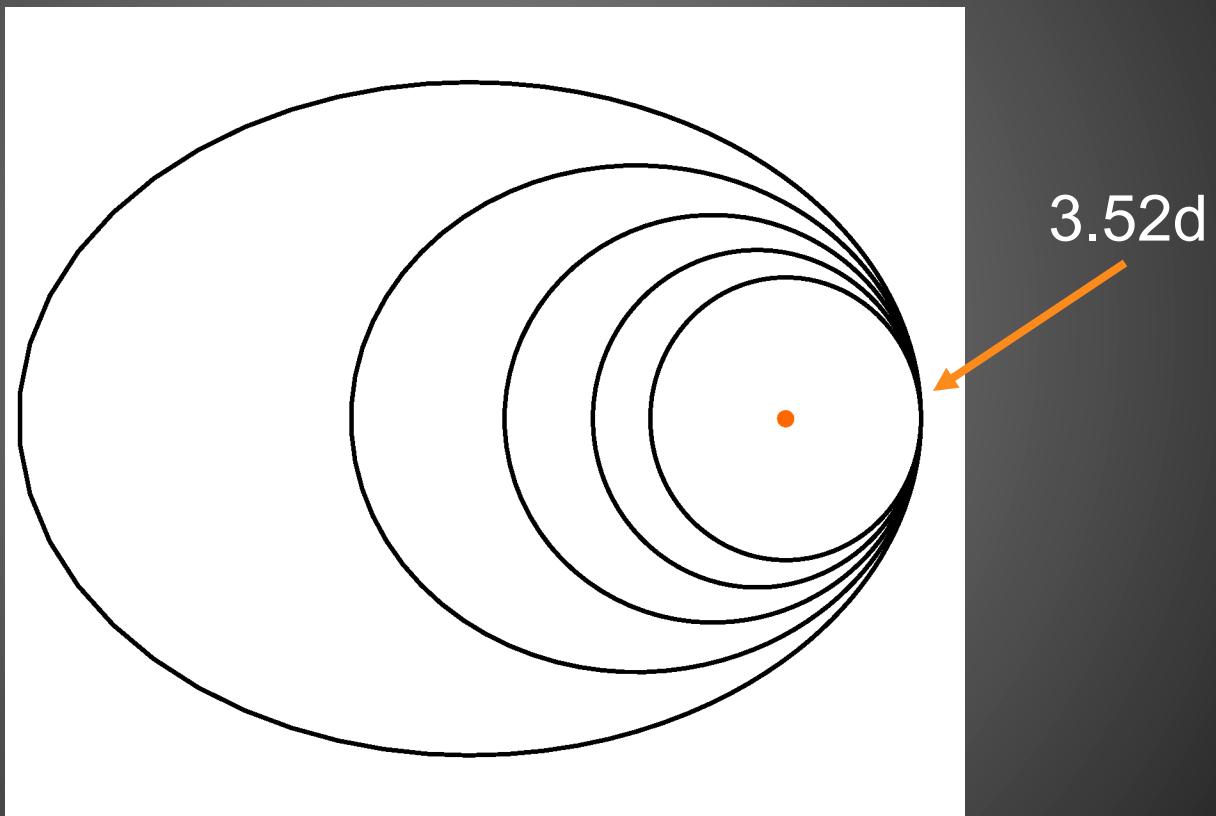


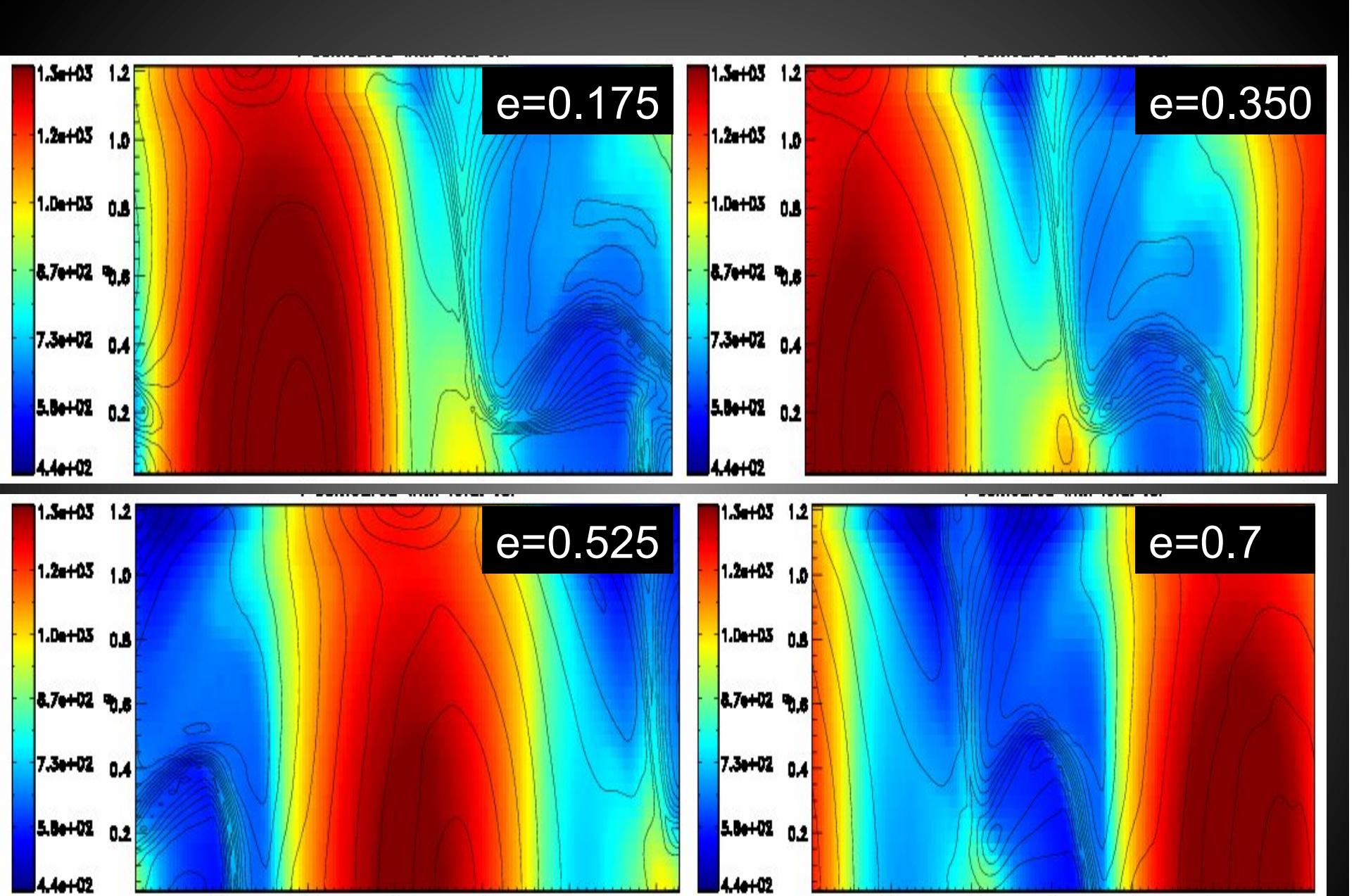
T_{peri}



$T_{peri} + 15\text{hr}$

Eccentric Planets





Conclusions

- Numerical treatment of radiation and dynamics must be included as coupled model.
- Three (pseudo) jets (one equatorial and two mid-lat.) are common features, with width decreasing with increased planetary rotation. Dissipation mechanisms are an open question.
- Optical and IR opacities both are important in determining location of stellar energy deposition and efficiency of redistribution to the night-side.
- Changing viscosity drastically alters streamlines, changing overall thermal structure
- Caution must be exercised in regions where material passes through strong shocks. Radial velocity becomes very important
- Dynamically driven variability causes variations transit spectra, but variation in hemispherically averaged phase curves will be difficult.
- Differences between ingress and egress transmission spectra may prove to be powerful tools for model diagnostics.
- Continuing observational programs, and coupling of dynamical and spectral models should allow tighter constraints on dynamical processes: eccentric planets, multiple wavelength (and continuous) observations, lower masses, younger planets.